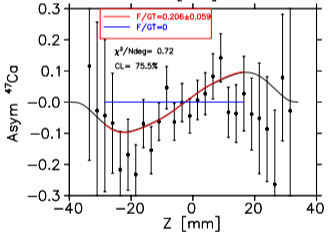
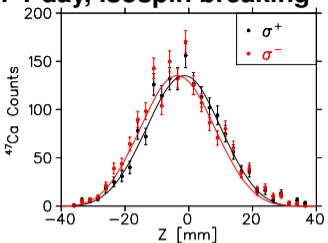
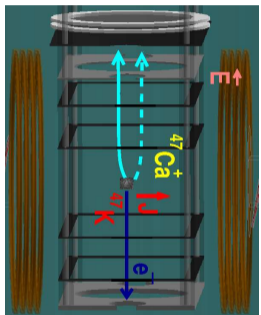


1000 atoms trapped for 1 day, isospin breaking ⁴⁷K Preliminary without weak mag



● Nonzero ⁴⁷Ca asymmetry wrt spin ⇒

a nonzero M_{Fermi}

$$M_F / M_{GT} =$$

$$0.21 \pm 0.06 \text{ stat} \pm ? \text{ syst} \Rightarrow$$

$$\langle \bar{A} | V_{Coulomb} | A \rangle =$$

$$160 \pm 50 \text{ stat} \pm ? \text{ syst keV}$$

● A_{recoil} is damped at extreme Z by a $\sim 6\%$ bkg from untrapped ⁴⁷K, measured by dedicated 'poof' tests

● Apparatus is symmetric:
X projection flat at 1σ to 0.05;
Unpolarized data has X, Z projections flat ~ 0.01

● β' 's fire the eMCP with $\sim 20\%$ quantum efficiency– these we measure to be ~ 0.002 correction

weak magnetism $2s_{1/2}$ nucleon

$$A_\beta \propto 1 + 0.005 E_\beta / E_0;$$

weak magnetism larger and d_i finite for the $20\% 1/2^+ \rightarrow 3/2^+$ transition

If $\langle \bar{A} | V_C | A \rangle = 160 \pm 50$ keV holds in ^{47}Ca , that's a

large fraction of $\langle \bar{A} | V_C | A \rangle = 0.35 \frac{\sqrt{n_1 n_2}}{2T} \frac{Z}{A^{2/3}} \text{MeV} = 190$ keV Auerbach Loc NPA2022
(unlike ^{56}Co 3 keV Markey, Boehm 1977 and ^{71}At 28 keV Severijns 2005)

\Rightarrow that single $1/2^+$ final state is the antianalog 😊, so the schematic ψ might be accurate Auerbach Loc NPA2022:

$$|A\rangle = \frac{1}{\sqrt{2T}} \left[\sqrt{n_1} \left| j_1^{n_1-1}(n) j_1(p) j_2^{n_2}(n) \right\rangle + \sqrt{n_2} \left| j_1^{n_1}(n) j_2^{n_2-1}(n) j_2(p) \right\rangle \right]$$

$$|\bar{A}\rangle = \frac{1}{\sqrt{2T}} \left[\sqrt{n_2} \left| j_1^{n_1-1}(n) j_1(p) j_2^{n_2}(n) \right\rangle - \sqrt{n_1} \left| j_1^{n_1}(n) j_2^{n_2-1}(n) j_2(p) \right\rangle \right]$$

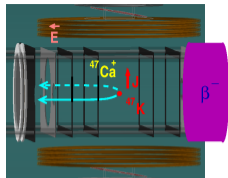
By inspection, the isovector piece of Herzceg 1965 P-even T-odd N-N interaction will flip the 'orthogonalizing' minus sign, enabling large matrix elements Barroso Blin-Stoyle 1973

$$V_{\text{t.v.}} = G_{\text{t.v.}} \frac{1}{2} [f(r) \hat{r} \cdot \mathbf{p} + \text{h.c.}] \times [1 + a \sigma^{(1)} \cdot \sigma^{(2)} (\tau_3^{(1)} + \tau_3^{(2)}) + (b + c \sigma^{(1)} \cdot \sigma^{(2)}) \tau_3^{(1)} \tau_3^{(2)}]$$

One still needs accurate $\psi(r)$ because $\hat{r} \cdot \mathbf{p}$ needs good tails! (Is this why Calaprice, Freedman never extracted microscopic TRV physics a, b, c from ^{56}Co ?)

So our ^{47}K goal: measuring isospin-enhanced TRV in a system that can be understood theoretically well enough to extract useful microscopic physics

For future TRV D:



In ^{56}Co , $E_1 = -0.01 \pm 0.02$,

$\langle \bar{A} | V_C | A \rangle = 2.9 \pm 0.5$,

$\langle \bar{A} | V_{TRV} | A \rangle = 54 \pm 110$ eV,

$M_{GT} = 0.0034$.

$\langle \bar{A} | V_C | A \rangle$ cancels in D or E : sensitivity scales with $1/M_{GT}$ ($=1/0.3$ in ^{47}K).

Measuring D to 0.001 in ^{47}K

(~ 3 weeks) leaves us 3x short in sensitivity to

$\langle \bar{A} | V_{TRV} | A \rangle$ compared to ^{56}Co , but the simpler TRV

N-N matrix elements are likely larger and calculable

^{47}K recoil order estimates still in progress

$^{47}_{19}\text{K}^{28}$ $\mu = 1.9 \mu_{\text{nucleon}} \Rightarrow$ thought to be 71% $2s_{1/2}$ Choudhary, Kumar, Srivasta, Suzuki PRC 103 064325 (2021)

Assuming $1/2^+ \rightarrow 1/2^+$ transition is $2s_{1/2} \rightarrow 2s_{1/2}$ (no orbital l contributions):

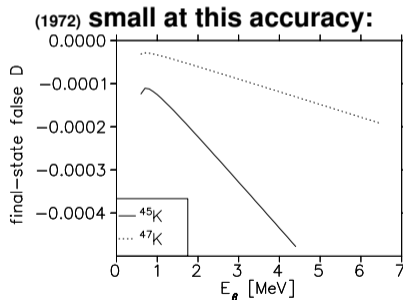
- Weak magnetism $b_W \sim$ the nucleon value
- 1st-class induced tensor $d_I \sim 0$

For our M_F/M_{GT} measurement,

A_β changed by ≤ 0.01

Recoil-order effects small at present level of accuracy \rightarrow statistics-limited measurement

Future D final-state effects Holstein PRC 5 1529



Note: ^{56}Co final-state $E_1=0.0002$ Calaprice 1977