

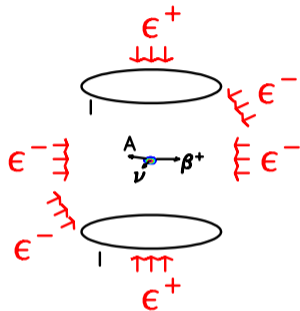
Discrete symmetries in β decay

- Parity P symmetry
 - How to test P symmetry experimentally
 - Only left-handed ν so far
- There is time-reversal symmetry violation
 \mathcal{T} in nature

There should? be more: 'baryogenesis'

\mathcal{P} , \mathcal{T} experiments at/with TRIUMF

- TRIUMF Neutral Atom trap (me 😊) β decay
 - Laser-polarized beamline at TRIUMF
- [Non-decays: • ν mixing,
- \mathcal{P} in Fr atoms; Searches for electric dipole moments (Kalita);
 - Using ultra-cold neutrons;
 - hopefully atomic fountain of francium atoms]



Symmetries: Continuous, Discrete

- Noether's theorem (1915):

Continuous symmetry	→	Conserved quantity
Time-translational invariance	→	Energy
Space-translational invariance	→	Momentum
Rotational invariance	→	Angular momentum
(Laplace-Runge-Lenz vector)	→	name?

THE LATE EMMY NOETHER.

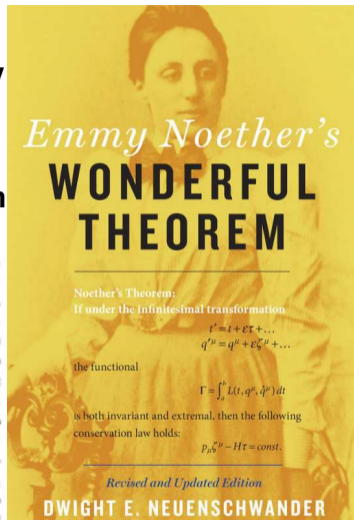
Professor Einstein Writes in Appreciation of a Fellow-Mathematician.

To the Editor of *The New York Times*:

In Ted Chiang's "Story of Your Life" [Movie "Arrival"]: aliens think in terms of the action, not position and momentum

gan. In the realm of algebra, in which the most gifted mathematicians have been busy for centuries, she discovered methods which have proved of enormous importance in the development of the present-day younger generation of mathematicians. Pure mathematics is, in its way, the poetry of logical ideas. One seeks the most general ideas of operation which will bring together in simple, logical and unified form the largest possible circle of formal relationships. In this effort toward logical beauty spiritual formulae are discovered necessary for the deeper penetration into the laws of nature.

- Discrete symmetries in quantum mechanics: Parity, Time reversal →



Historical Ideas about P , T breaking

- Wigner considered implications of P , T symmetry conservation in atomic spectra 1926-28. Showed $\langle T\psi_i, T\psi_f \rangle = \langle \psi_f, \psi_i \rangle^*$

“In quantum theory, invariance principles permit even further reaching conclusions than in classical mechanics.” (D. Gross, Physics Today 48 46 (1995))

- Weyl 1931 considered C , P , T and CPT in “Maxwell-Dirac theory”: $C \Rightarrow$ Dirac eq. negative energy states had to have same mass as the e^- plato.stanford.edu

- From “CP Violation Without Strangeness” Khriplovich and Lamoreaux: 1949 Dirac “I do not believe there is any need for physical laws to be invariant under reflections in space and time although the exact laws of nature so far known do have this invariance.”

- 1956 Lee and Yang proposed \not{P} in weak decays to fix the θ - τ puzzle

- Feynman gives Ramsey 50:1 odds \not{P} would not be observable
Ramsey experiment starting at ORNL gets derailed by fission experiments...
it's OK, Ramsey won 1989 Nobel for his fringes

- 1957 3 simultaneous experimental measurements of $\not{P} \rightarrow$

Parity (From A. Zee “Fearful Symmetry”)

As of 1956, we thought
all interactions
respected parity

Parity operator

$$P \psi(\vec{r}) \rightarrow \pm \psi(-\vec{r})$$

P



=



1957:

$\tau - \theta$ Puzzle

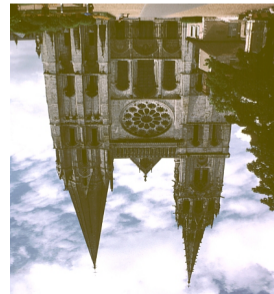
+ μ decay

+ ^{60}Co decay \Rightarrow

P



\neq



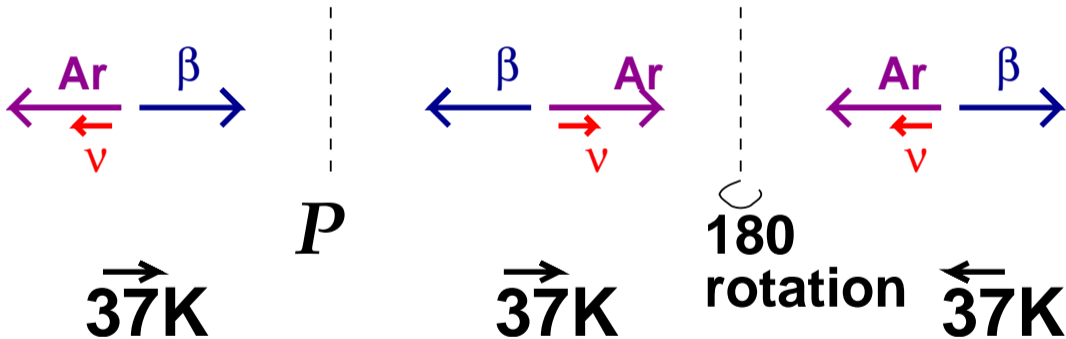
Decays: Parity Operation can be simulated by Spin Flip

Under Parity operation P :

$$\vec{r} \rightarrow -\vec{r}$$

$$\vec{p} \sim \frac{d\vec{r}}{dt} \rightarrow -\vec{p}$$

$$\vec{J} = \vec{r} \times \vec{p} \rightarrow +\vec{J}$$



\Rightarrow A spin flip corresponds exactly to P reversal

Decays don't exactly test T -reversal symmetry



One experimental discovery of parity violation

Wu, Ambler, Hayward,
Hopper, Hobson,
PR 105 1413 Feb '57
Dilution Refrigerator to
spin-polarize



$$W[\theta] = 1 + PA\hat{\mathbf{I}} \cdot \frac{\vec{p}_\beta}{E_\beta}$$

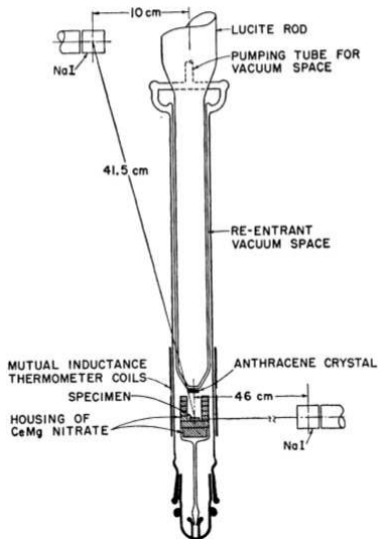
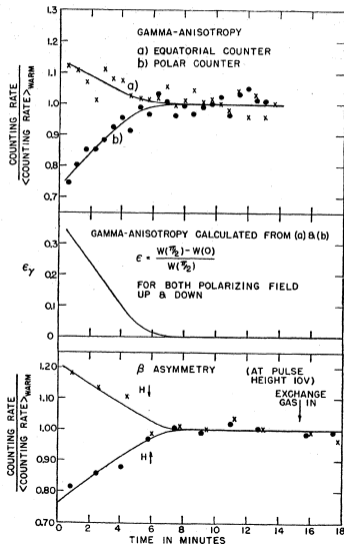
$$= 1 + A\frac{v}{c}\cos[\theta]$$

$$A_{\beta^-} \approx -1.0$$

Followup:
 ${}^{58}\text{Co} \rightarrow {}^{58}\text{Fe} + \beta^+ + \nu$

$$A_{\beta^+} > 0$$

CP conserved?



Measure ν helicity: transfer to γ circular polarization

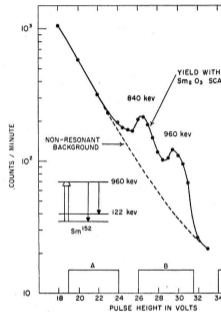
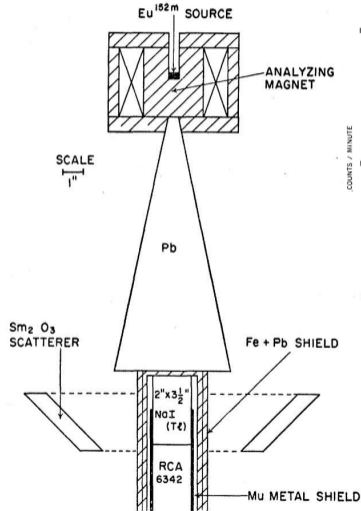
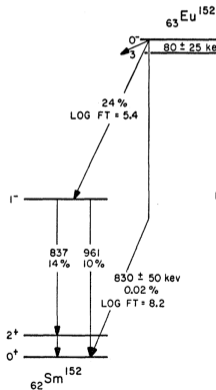
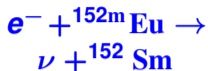
Goldhaber, Grodzins, Sunyar
Phys Rev 109 1015 (Dec 1957)

- Upward-going ν populates $\langle I_z \rangle = 0, +1$ **not -1**

- So γ is circularly polarized—transmission through magnet depends on iron polarization:
 $\frac{N_+ - N_-}{N_+ + N_-} = 0.017 \pm 0.003$

- Upward ν boosts γ momentum so it can be absorbed on-resonance
 $\Rightarrow \nu$ helicity $-1 \pm 10\%$

- ($\bar{\nu}$ helicity $\sim +1$)
Palathingal PRL 524 24 '69)



TRIUMF Physics and time reversal

When $t \rightarrow -t$, does anything change?

- Wave Equ. is 2nd-order in t : $\nabla^2 u = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2}$ **symmetric in t**

- Heat Equ. is 1st-order in t : $\nabla^2 u = -\frac{\partial u}{\partial t}$ **$t \rightarrow -t$, boom?**

‘Dissipation’, like friction... The arrow of time remains a research problem in stat mech, but it’s not from (known) particle physics

- Schroedinger Equation is 1st order: $i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2}$

‘Take the complex conjugate’

(but see Dressel et al. PRL 119 220507 (2017)

“Arrow of Time for Continuous Quantum Measurements”)

Microscopic physics was thought to be symmetric in t

Simulating \mathcal{T} in decays?

We've constructed an angular correlation, a scalar observable, by a dot product of two vectors

$$1 + \hat{p} \cdot \hat{J}$$

which is odd under P as we need

(\vec{p} is even, $\vec{J} = \vec{r} \times \vec{p}$ is odd)

But \vec{J} is odd under T , not even

So we need at least 3 vectors to have a T-odd scalar observable,

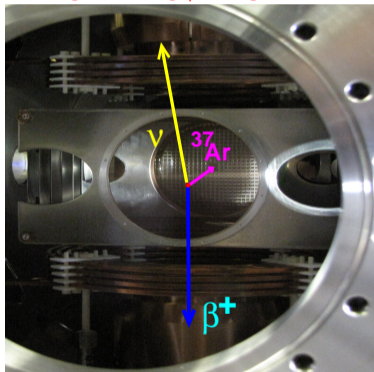
the scalar triple product $\vec{v}_1 \cdot (\vec{v}_2 \times \vec{v}_3)$

An example \rightarrow

TRIUMF \mathcal{T} correlation of 3 of 4 momenta

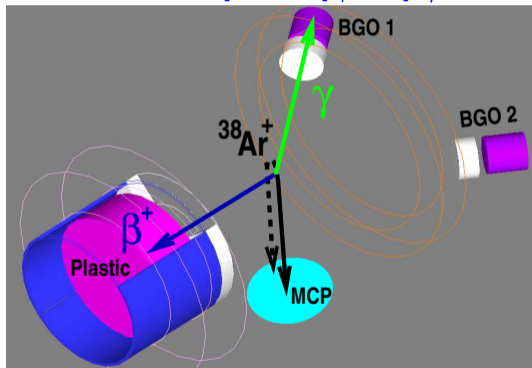
$$\mathbf{t} \rightarrow -\mathbf{t} \Rightarrow \vec{\mathbf{p}} \propto \frac{d\vec{\mathbf{r}}}{dt} \rightarrow -\vec{\mathbf{p}}$$

$$\text{but } \vec{\mathbf{p}}_{\text{recoil}} \cdot \vec{\mathbf{p}}_{\beta} \times \vec{\mathbf{p}}_{\nu} \equiv 0 \text{ ☹}$$



$$\vec{\mathbf{p}}_{\nu} \cdot \vec{\mathbf{p}}_{\beta} \times \vec{\mathbf{p}}_{\gamma} = -\vec{\mathbf{p}}_{\text{recoil}} \cdot \vec{\mathbf{p}}_{\beta} \times \vec{\mathbf{p}}_{\gamma}$$

$$\xrightarrow{\mathbf{t} \rightarrow -\mathbf{t}} \vec{\mathbf{p}}_{\text{recoil}} \cdot \vec{\mathbf{p}}_{\beta} \times \vec{\mathbf{p}}_{\gamma}$$



- We can test symmetry of apparatus with coincident pairs ☺
- Not exact. Outgoing particles interact \rightarrow fake \mathcal{T}

EDM in a fundamental particle breaks T : this is exact

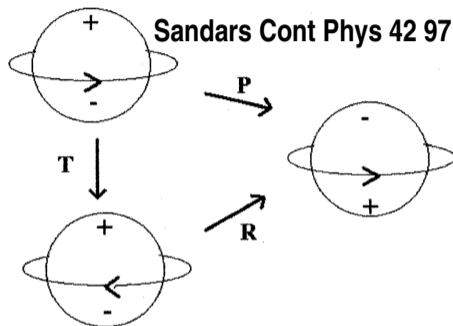
Landau, Nucl. Phys. 3 (1957) p. 127

Electric Dipole moment $\vec{d} = \sum q_i \vec{r}_i$

Since the angular momentum is the only vector in the problem, $\vec{d} = a\vec{J}$

Under T , $\vec{J} \xrightarrow{T} -\vec{J}$ $\vec{d} \xrightarrow{T} +\vec{d}$

If the physics is invariant under T , this is a contradiction, $\Rightarrow a = 0$



[• The other logical possibility: there are 2 states, with opposite sign of the EDM, and T just formally changes one state to the other.

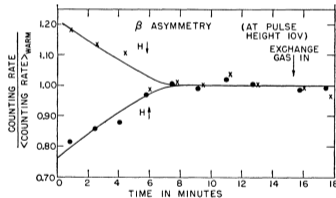
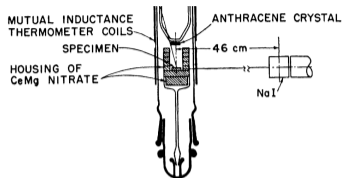
For most fundamental particles, we know there aren't 2 states

Why do we know the electron doesn't have 2 states?

E.g. some polar molecules have a dipole moment listed in tables, which produces degenerate states and does not break T ...]



Parity broken, why not Time?



Immediately after \mathcal{P} arity was seen to be totally broken in β decay (' ν left-handed')

**Wu, Ambler, Hayward, Hopper, Hobson,
PR 105 (1957) 1413**

Many T-odd observables were proposed:

PHYSICAL REVIEW

VOLUME 106, NUMBER 3

Possible Tests of Time Reversal Invariance in Beta Decay

J. D. JACKSON,* S. B. TREIMAN, AND H. W. WYLD, JR.

Palmer Physical Laboratory, Princeton University, Princeton, New Jersey

(Received January 28, 1957)

Need scalar triple products of 3 vectors:
observables involving spin

$$D \hat{J} \cdot \frac{\vec{p}_\beta}{E_\beta} \times \frac{\vec{p}_\nu}{E_\beta} \quad R \vec{\sigma}_\beta \cdot \hat{J} \times \frac{\vec{p}_\beta}{E_\beta} \quad \text{TRINAT has } D \text{ ideas}$$

are consistent with $\mathcal{T} < 0.001$

(We're looking for \mathcal{T} that could still be big:) but some has been found \rightarrow

Possible Tests of Time Reversal Invariance in Beta Decay

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~~CP~~ discovered in $K\bar{K}$ meson decays in 1963,
though not much (Cronin and Fitch Nobel prize 1980)

Quark eigenstates in the weak interaction:

To explain some weak decays,

$$|u\rangle \rightarrow |d\rangle + \epsilon|s\rangle \quad \text{i.e.} \quad |u\rangle \rightarrow \cos(\theta_C)|d\rangle + \sin(\theta_C)|s\rangle$$

For 3 families of particles,

→ 3x3 unitary “CKM” matrix between $|d\rangle$, $|s\rangle$, $|b\rangle$

There is one complex phase, which leads to this type of ~~CP~~

A reason for 3 generations of particles?

That one phase is consistent with CP in $K\bar{K}$ and $B\bar{B}$ systems

There have been hints in $K\bar{K}$ and $B\bar{B}$ of more CP than in the standard model,

$p\bar{p} \rightarrow \mu^+\mu^+$ or $\mu^-\mu^-$ CP at 3.6σ Abazov PRD 2014

Fermilab;
so this 2001 cartoon was a little premature \rightarrow



J. Faberge. CERN Courier, 6, No. 10, 193 (October 1966). [Courtesy of Madame Faberge.]

T2K ν_μ oscillations different from $\bar{\nu}_\mu$ at 2 to 3 σ Nature 580 339 (2020)

CP could have some utility for cosmology \rightarrow

The excess of matter over antimatter can come from \mathcal{CP}

Sakharov JETP Lett 5 24 (1967) used \mathcal{CP} to generate the universe's excess of matter over antimatter:

- \mathcal{CP} ,
- baryon nonconservation, and
- nonequilibrium.

But known \mathcal{CP} is too small by 10^{10} , so 'we' need more to exist. Caveats:

- You could use \mathcal{CPT} [Dolgov Phys Rep 222 (1992) 309]
- We need \mathcal{CP} in the early universe, not necessarily now

So we look for more \mathcal{CP} . How is this related to \mathcal{T} ?

**this seems a
 little abstract
 concrete
 demonstrative
 example from
 Ramsey-Musolf
 at INT this week**
CP explaining
 T2K's ν vs. $\bar{\nu}$
 result lets heavy
N decay this
 way in some
 models

Neutrinos and the Origin of Matter

- *Heavy neutrinos decay out of equilibrium in early universe*
- *Majorana neutrinos can decay to particles and antiparticles*
- *Rates can be slightly different (CP violation)*

$$\Gamma(N \rightarrow \ell H) \neq \Gamma(N \rightarrow \bar{\ell} H^*)$$

- *Resulting excess of leptons over anti-leptons partially converted into excess of quarks over anti-quarks by Standard Model sphalerons*

\mathcal{T} is related to \mathcal{CP} by the “CPT Theorem”

“All local Lorentz invariant QFT’s are invariant under CPT”

Schwinger Phys Rev 82 914 (1951)

Lüders, Pauli, Bell 1954

- Gravity \rightarrow not flat:

K meson experiments Adler PhysLettB 364 (1995) 239 test

\mathcal{CPT} to within 1000x expected from quantum gravity

- Strings not ‘local’

Proofs still pursued \rightarrow

Assuming CPT, $\mathcal{CP} \Leftrightarrow \mathcal{T}$ in most physics theories

The matter excess then motivates \mathcal{T} searches

Studies in History and Philosophy of Modern Physics 45 (2014) 46–65



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On the CPT theorem

Hilary Greaves^{a,*}, Teruji Thomas^{b,1}

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CPT theorem
Discrete symmetries
Spacetime symmetries

ABSTRACT

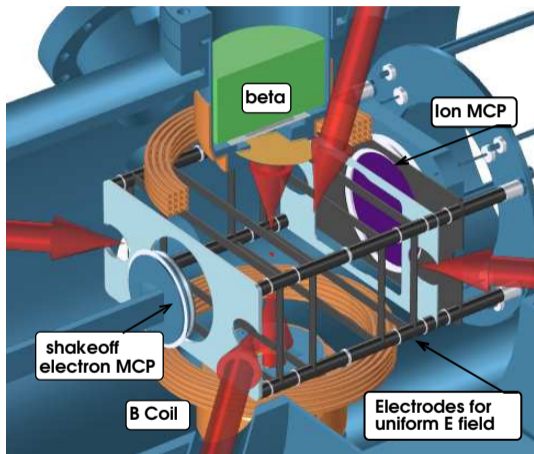
We provide a careful development and rigorous proof of the CPT theorem within the framework of mainstream (Lagrangian) quantum field theory. This is in contrast to the usual rigorous proofs in purely axiomatic frameworks, and non-rigorous proof-sketches in the mainstream approach. We construct the CPT transformation for a general field directly, without appealing to the enumerative classification of representations, and in a manner that is clearly related to the requirements of our proof. Our approach applies equally in Minkowski spacetimes of any dimension at least three, and is in principle neutral between classical and quantum field theories: the quantum CPT theorem has a natural classical analogue. The key mathematical tool is that of complexification; this tool is central to the existing axiomatic proofs, but plays no overt role in the usual mainstream approaches to CPT.

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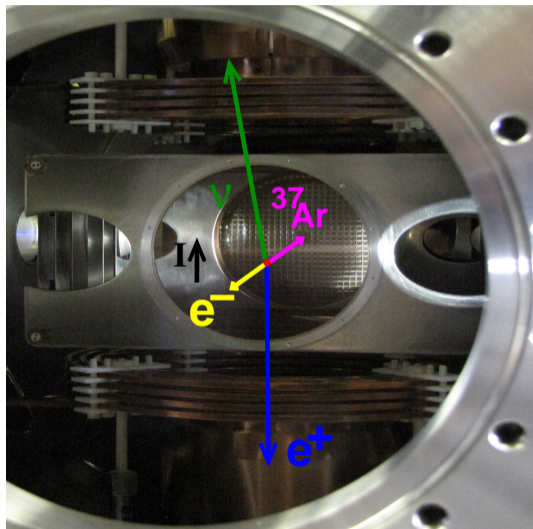


^{37}K decay geometry

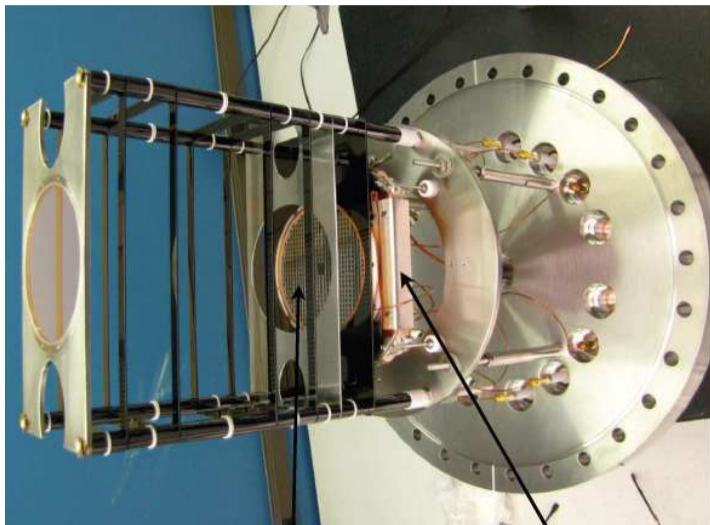


- β , recoil nucleus
- shakeoff e^- for TOF trigger

The decay pattern shown on the right is helicity-forbidden if the ν goes straight up



ion MCP assembly



14 inch CF flange

Electrostatic field

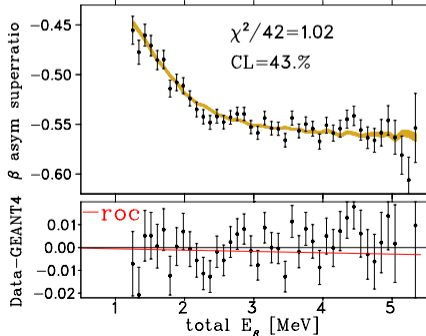
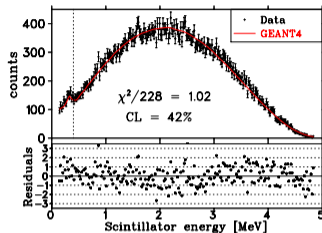
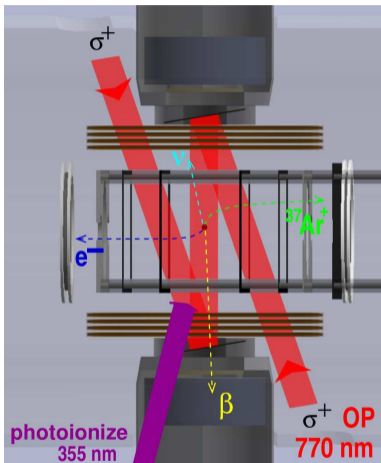
**delay-line anode for
position info**

no stray wires

**Low-Z (glassy carbon,
titanium) to minimize β^+
scattering**



β^+ asymmetry ^{37}K data



Fenker et al. Phys Rev Lett 120, 062502 (2018)

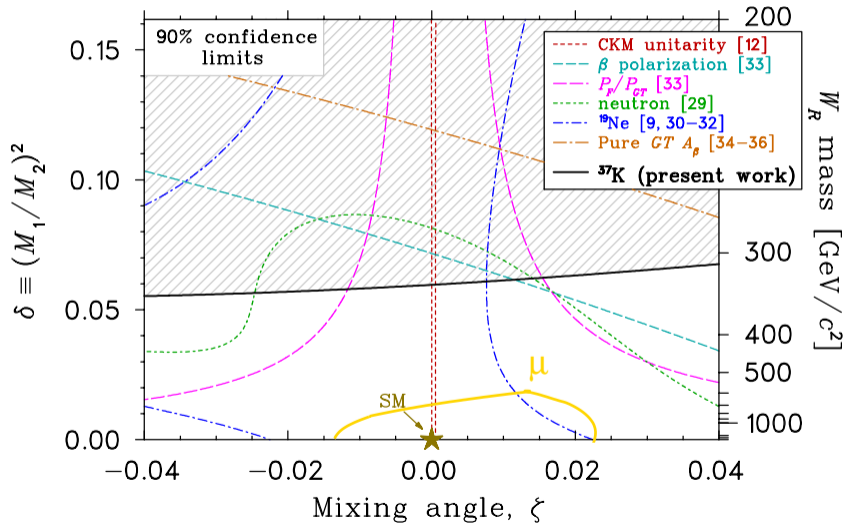
A_β [experiment] =
 -0.5707 ± 0.0019

A_β [theory] =
 -0.5706 ± 0.0007

The best fractional accuracy achieved in nuclear or neutron β decay



Still no wrong-handed ν 's



Extra W' with heavier mass, couples to wrong-handed ν_R

We can evade TWIST limits by assuming the muon ν_R is heavy
LHC $M'_W > 3.7$ TeV
90%



The nucleon: a special place for γ 's

S.M. interactions combined in the nucleon:

Harvey Hill Hill PRL 2007 Gardner He PRD 2013

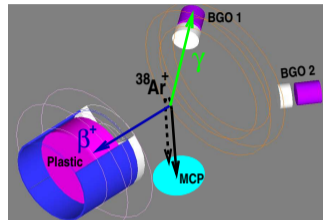
$$\mathcal{L} = \frac{-4c_5}{m_{\text{nucleon}}^2} \frac{eG_F V_{ud}}{\sqrt{2}} \epsilon^{\sigma\mu\nu\rho} \overset{\text{QCD}}{\downarrow} \bar{p} \gamma_\sigma n \overset{\text{Weak decay}}{\downarrow} \psi_e L \gamma_\mu \psi_\nu \overset{\text{E\&M}}{\downarrow} F_{\nu\rho}$$

Interference with S.M. β decay 'vector current' gives $\beta\nu\gamma$ decay contribution with the scalar triple product we want:

$$|\mathcal{M}_{c_5}|^2 \propto \frac{\text{Im}(c_5 g_V)}{M^2} \frac{E_e}{p_e k} (\vec{p}_e \times \vec{k}_\gamma) \cdot \vec{p}_\nu$$

- This source of \mathcal{T} scales with p_{lepton}^2 , so is $\sim 10^2$ larger in ^{37}K decay than neutron

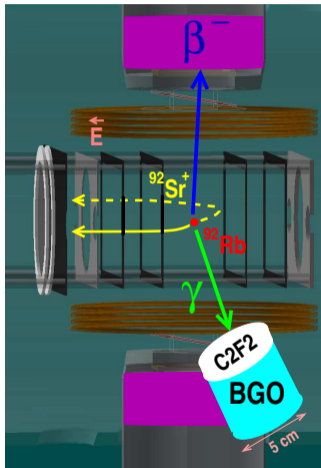
- Direct constraint from $n \rightarrow p \beta\nu\gamma$ branch $\propto |c_5|^2$
Bales PRL 2016: $3.4 \pm 0.2 \times 10^{-3}$ (theory 3.1×10^{-3})
 $\Rightarrow \frac{\text{Im}(c_5)}{M^2} \leq 8 \text{MeV}^{-2} \Rightarrow ^{37}\text{K} \mathcal{T}$ asym can still be ~ 1 😊



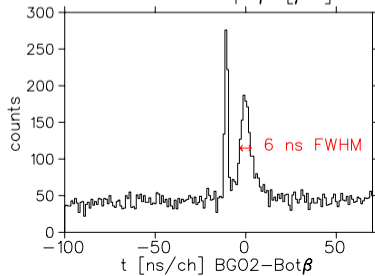
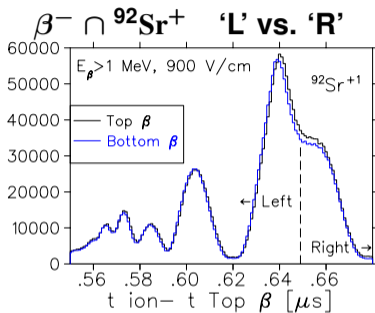
\mathcal{T} needs new physics with scale $M \sim \text{MeV}$



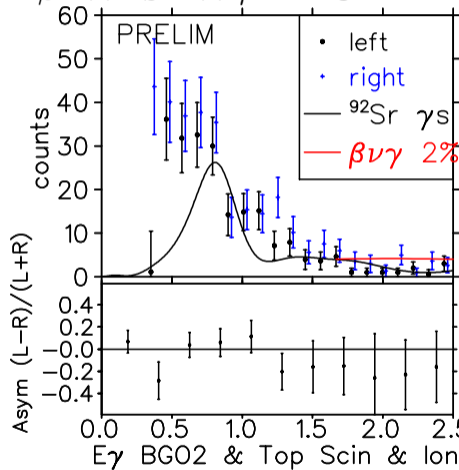
\mathcal{T} test with $^{92}\text{Rb } 0^- \rightarrow ^{92}\text{Sr } 0^+ + \beta^- \nu \gamma$



BGO \rightarrow GAGG
(Ce:Gd₃Al₂Ga₃O₁₂)



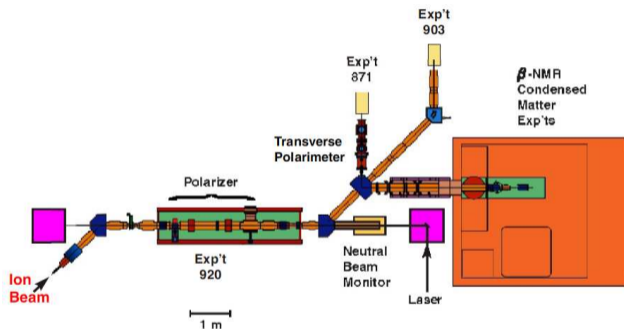
$\beta^- \cap ^{92}\text{Sr}^+ \cap \gamma$ 'L' vs. 'R'



Sensitivity to
 $\lesssim 2\% \beta\nu\gamma$ branch

TRIUMF Laser-Polarized beam at TRIUMF/ISAC

C.D.P. Levy et al. / Nuclear Physics A 746 (2004) 206c–209c



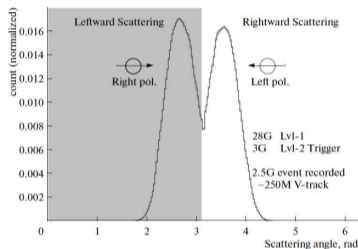
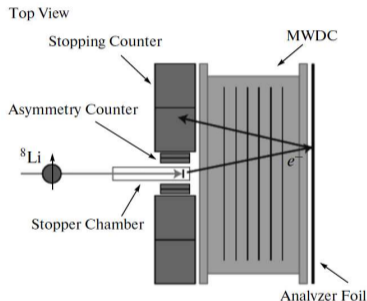
- 50-70% polarization, 20-50% efficient
- Re-stripped +1 beam deliverable to several beamlines

- Used for aligned ^{20}Na β correlation (2nd-class current comparison with ^{20}F) K. Minamisono PRC 84 055501 (2011)

^8Li R, Jiro Murata, Rikkyo U.

TRV possibilities include ^{36}K E and ^{20}Na β -delayed α energy shift (Clifford PRL 50 (1983) 23)

Mott scattering Time reversal Violation progress



$$R\vec{\sigma}_\beta \cdot \hat{\mathbf{J}} \times \vec{v}_\beta \xrightarrow{t \rightarrow -t} \rightarrow$$

$$-R\vec{\sigma}_\beta \cdot \hat{\mathbf{J}} \times \vec{v}_\beta$$

**Totsuka et al Phys Part
Nuclei 45 244 (2014)
Small false asymmetry in
rectangular geometry**

→ a more symmetric cylindrical geometry, finished data-taking Dec 2017

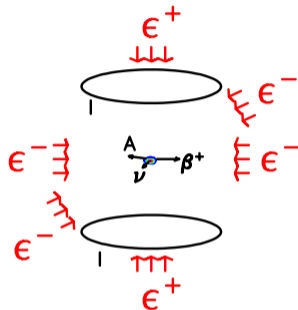
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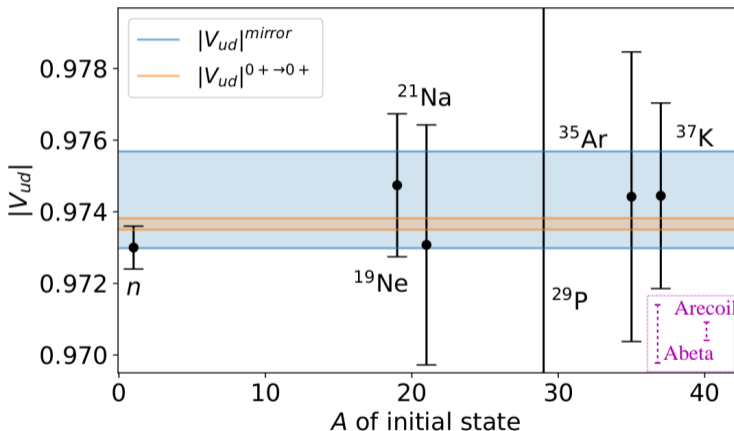
\mathcal{P} , \mathcal{T} experiments at/with TRIUMF

- TRIUMF Neutral Atom trap (me 😊) β decay
 - Laser-polarized beamline at TRIUMF
- [Non-decays: • ν mixing,
- \mathcal{P} in Fr atoms; Searches for electric dipole moments (Kalita);
 - Using ultra-cold neutrons;
 - hopefully atomic fountain of francium atoms]





Weak interaction: same strength, all nuclei?



Deduced V_{ud}
from mirror decays

Are people overestimating
their uncertainties? We
aren't 😊

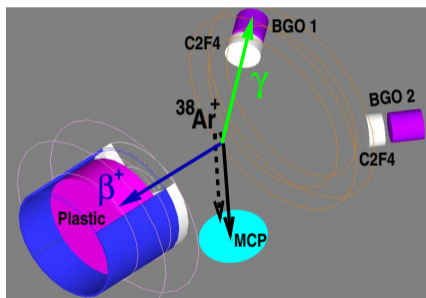
We project to reach 0.0005
accuracy, as good as any
 $0^+ \rightarrow 0^+$ except ^{26m}Al .

Assumes 5% isospin
breaking calculation.

Hayen and Severijns, arXiv:1906.09870 (June 2019)

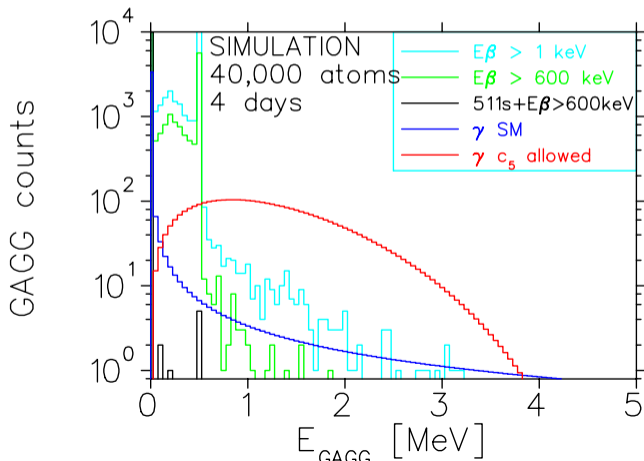
Simulations: E_γ signature and backgrounds

- Classical bremsstrahlung $\propto 1/E_\gamma$
- Any time-reversal violating interaction involves β, ν and $\gamma \Rightarrow$ 4-body phase space $\propto E_\gamma(Q - E_\gamma)^3$ Bernard PLB 593 (2004)

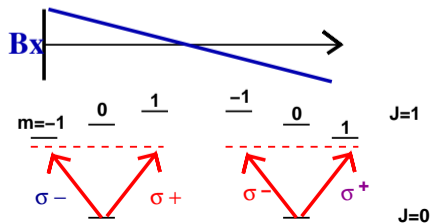
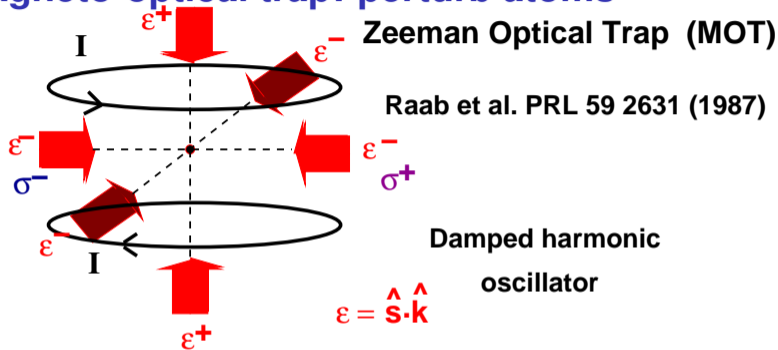


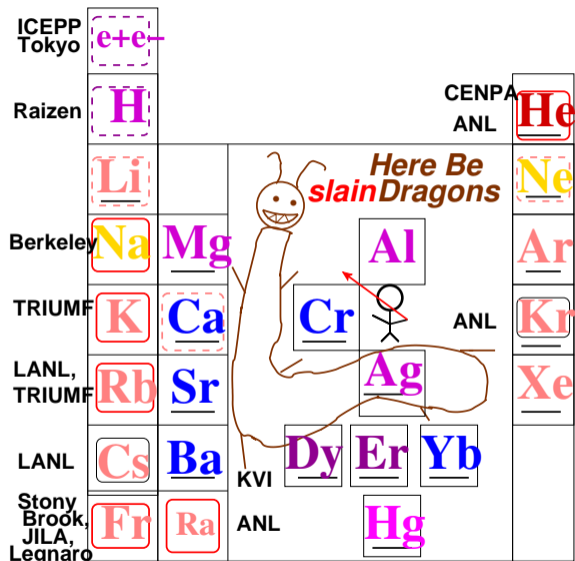
We are concentrating on:

- $E_\gamma > 511$ keV
- the β^+ in the opposite detector



Magneto-optical trap: perturb atoms

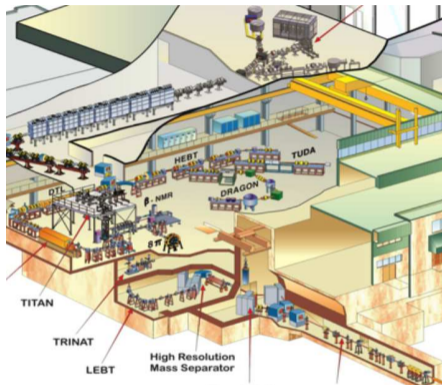




What elements can be laser cooled?



TRiumf Neutral Atom trap at ISAC

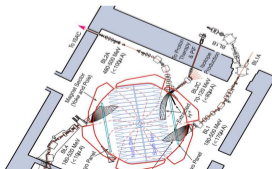


^{37}K $8 \times 10^7/\text{s}$

TiC target
 1750°C

$70 \mu\text{A}$
protons

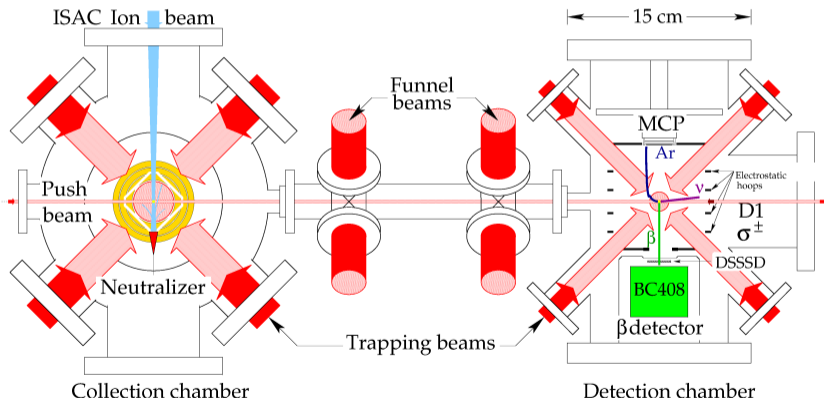
main TRIUMF cyclotron
'world's largest'
 500 MeV H^- (0.5 Tesla)





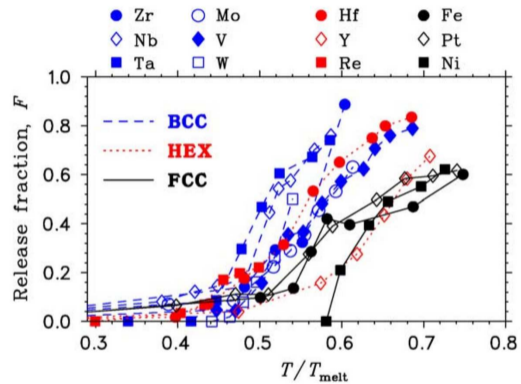
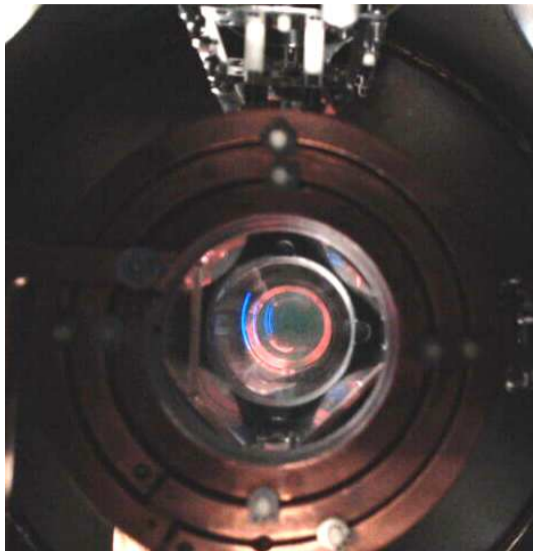
TRINAT plan view

- Isotope/Isomer selective
- Avoid untrapped atom background with 2nd trap
- 75% transfer
- 0.7 mm cloud for β -Ar⁺ \rightarrow ν momentum



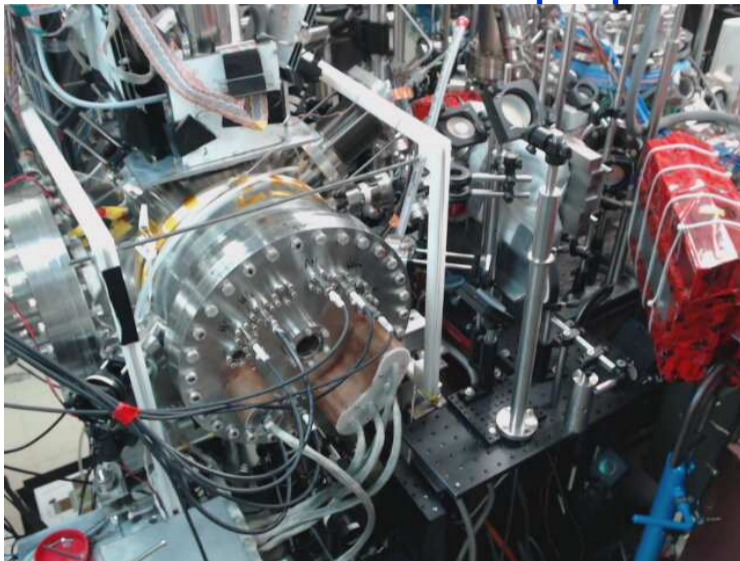
- Spin-polarized $99.1 \pm 0.1\%$

Neutralizer and Collection trap

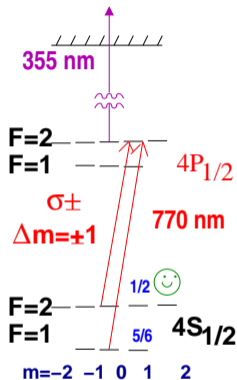




TRINAT lab: “tabletop experiment”



TRIUMF Optical pumping and probing ^{37}K



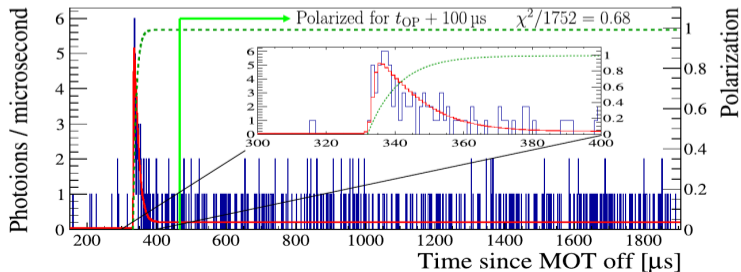
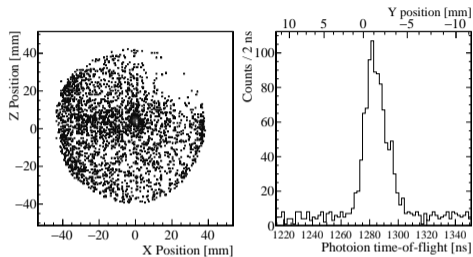
Photoionize 1%

in situ probe

$P_{+} = +0.9913(8)$

$P_{-} = -0.9912(9)$

Fenker NJP 2016



Preview: Weak interaction breaks parity: Consequences?

'Pulsar kicks'



IGR J11014-6103

$v = 0.01 c$ →

Fuller PRD 2003

Forced $p + e^- \rightarrow n + \nu$

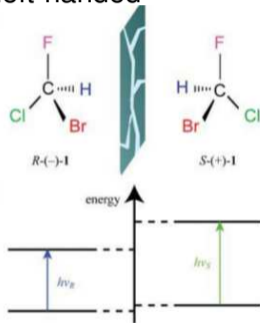
$W(\theta) = 1 + \frac{\langle m_\nu \rangle}{I} A_\nu \cos(\theta_i)$

B field polarizes p 's

Need ν_e to include 10^{-8}

admixture of $m_\nu \sim \text{keV}$

Earthling's amino acids are all left-handed



Letokhov PLA'75

Darquie CHIRALITY 2010

$\Delta E \sim 10^{14-16} \text{eV}$

Not Enough for left-handed

bugs to win, so →

Spin-polarized SN ν 's could preferentially zap wrong-handed amino acids

Finding the right environment for spin-polarized amino acids? e.g. :

Astrobiology 18 (2018)

Selection of Amino Acid

Chirality via ν Interactions

with ^{14}N in $\vec{E} \times \vec{B}$ Fields

M.A. Famiano, R.N. Boyd

(TRIUMF EEC 90's)...