mtv

🖻 in atoms

extras

Parity and time-reversal symmetry in β decay

- Parity *P* symmetry
 How to test *P* symmetry experimentally
 Only left-handed ν so far
- There is time-reversal symmetry violation *T* in nature

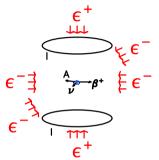
There should? be more: 'baryogenesis'

- TRIUMF Neutral Atom trap (me S) β decay
- Laser-polarized beamline at TRIUMF

[Non-decays: • ν mixing,

- *P* in Fr atoms; Searches for electric dipole moments;
- Using ultra-cold neutrons;
- hopefully atomic fountain of francium atoms]

https://trinat.triumf.ca/talks/triumf-summer-undergrad-talk-2021-jb



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extras

Symmetries: Continuous, Discrete

- Noether's theorem (1915):
 - Continuous symmetry Time-translational invariance

time

Space-translational invariance

Rotational invariance

(Laplace-Runge-Lenz vector)

THE LATE EMMY NOETHER.

Professor Einstein Writes in Appreolation of a Fellow-Mathematician.

To the Editor of The New York Times :

In Ted Chiang's "Story of Your Life" [Movie "Arrival"]: aliens think in terms of the action, not position and momentum \rightarrow Conserved quantity

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- Energy
- Momentum
- \rightarrow Angular momentum

name?

gan. In the realm of algebra, in which the most gifted mathematicians have been busy for centuries, she discovered methods which have proved of enormous importance in the development of the present-day younger generation of mathematicians. Pure mathematics is, in its way, the poetry of logical ideas. One socks the most general ideas of operation which will bring together in simple, logical and unified form the largest possible circle of formal relationships. In this effort toward logical beauty spiritual formulae are discovered necessary for the deeper penetration into the laws of nature.

Emmy Noether's WONDERFUL THEOREM

Nother's Theorem: If under the infinitesimal transformation $\begin{aligned}
\mu^{r} &= 4r + \xi \xi^{-} + \dots \\
q^{r} &= q^{r} + \xi \xi^{-} + \dots \\
\mu^{r} &= \int_{0}^{1} \mathcal{U}_{1} q^{r} dq^{r} dq^{r} dt
\end{aligned}$ the functional $\begin{aligned}
\mu &= \int_{0}^{1} \mathcal{U}_{1} q^{r} dq^{r} dt
\end{aligned}$ is both invariant and extremal, then the following conservation law holds: $\mu_{n} \xi^{r} - H \epsilon - const.$ Revised and Updated Edition **DWIGHT E. NEUENSCHWANDER**

 \bullet Discrete symmetries in quantum mechanics: Parity, Time reversal \rightarrow

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Historical Ideas about P, T breaking

• Wigner considered implications of P, T symmetry conservation in atomic spectra 1926-28. Showed $\langle T\psi_i, T\psi_f \rangle = \langle \psi_f, \psi_i \rangle^*$

"In quantum theory, invariance principles permit even further reaching conclusions than in classical mechanics." (D. Gross, Physics Today 48 46 (1995))

- Weyl 1931 considered *C*, *P*, *T* and *CPT* in "Maxwell-Dirac theory": $C \Rightarrow$ Dirac eq. negative energy states had to have same mass as the e^- plato.stanford.edu
- From "CP Violation Without Strangeness" Khriplovich and Lamoreaux: 1949 Dirac "I do not believe there is any need for physical laws to be invariant under reflections in space and time although the exact laws of nature so far known do have this invariance."
- 1956 Lee and Yang proposed ${\cal P}$ in weak decays to fix the θ - τ puzzle
- Feynman gives Ramsey 50:1 odds $\not P$ would not be observable Ramsey experiment starting at ORNL gets derailed by fission experiments... it's OK, Ramsey won 1989 Nobel for his fringes

parity

time

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Parity (From A. Zee "Fearful Symmetry")

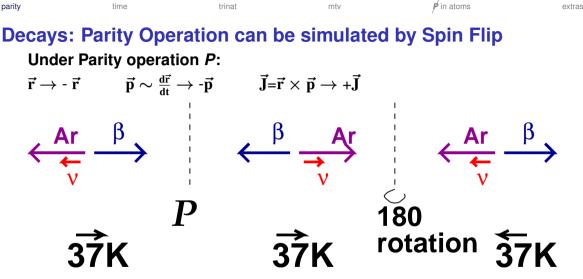
As of 1956, we thought all interactions respected parity Parity operator $P \psi(\vec{r}) \rightarrow \pm \psi(-\vec{r})$





1957: $\tau - \theta$ Puzzle + μ decay + 60 Co decay \Rightarrow





 \Rightarrow A spin flip corresponds exactly to *P* reversal Decays don't exactly test *T*-reversal symmetry parity

time

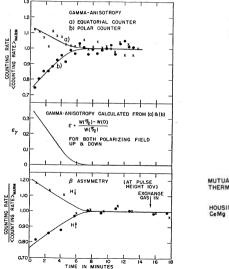
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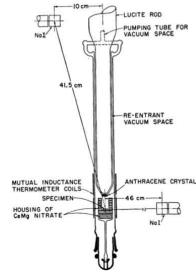
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R One experimental discovery of parity violation





Wu, Ambler, Hayward, Hopper, Hobson. PR 105 1413 Feb '57 **Dilution Refrigerator to** spin-polarize 60 Co $\rightarrow ^{60}$ Ni + β^- + $\bar{\nu}$ $W[\theta] = 1 + PA\hat{I} \cdot \frac{p_{\beta}}{F_{\alpha}}$ $= 1 + A \frac{v}{c} \cos[\theta]$ $A_{\beta-} \approx -1.0$ Followup:

 $A_{\beta+} > 0$ CP conserved?

 ${}^{58}\text{Co} \rightarrow {}^{\overline{58}}\text{Fe} + \beta^+ + \nu$

parity time trinat mtv p[≠] in atoms

Measure ν helicity $\epsilon = \hat{s}_{\nu} \cdot \hat{k}_{\nu}$ directly: transfer \hat{s}_{ν} to γ circular polarization; boost \vec{k}_{γ} by $\pm \vec{k}_{\nu}$

Goldhaber, Grodzins, Sunvar $e^- + {}^{152m} Eu \rightarrow$ Fuist SOURCE Phys Rev 109 1015 (Dec 1957) $\nu + {}^{152}$ Sm ANALYZING Smg Os SCA MAGNET • Upward-going ν populates 840 ke 63Eu $\langle I_z \rangle = 0, \pm 1 \text{ not -1}$ NON-RESONANT 80 1 25 • So γ is circularly polarized– SCALE 108 transmission through magnet 960 ke 24% LOG FT = 5.4 depends on iron polarization: $\frac{N_{+}-N_{-}}{N_{+}+N_{-}}=0.017\pm0.003$ Pb • Upward ν boosts γ Sma Ox momentum so it can be SCATTERER Fe + Ph SHIFLD absorbed on-resonance 837 961 $\Rightarrow \nu$ helicity -1 \pm 10% 830 ± 50 kev 0.02 % LOG FT = 8.2 (• $\bar{\nu}$ helicity \sim +1 Palathingal PRL 524 24 '69) RCA MU METAL SHIFLD 62</sub>Sm¹⁵²

Surprisingly enough, this is the best direct measurement of ν helicity = $\hat{s_{\nu}} \cdot \hat{k_{\nu}}$

RIUMF Physics and time reversal

time

parity

When t \rightarrow -t, does anything change?

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• Wave Equ. is 2nd-order in t: $\nabla^2 u = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2}$ symmetric in t

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p in atoms

- Heat Equ. is 1st-order in t: $\nabla^2 u = -\frac{\partial u}{\partial t}$ t \rightarrow -t, boom? 'Dissipation', like friction... The arrow of time remains a research problem in stat mech, but it's not from (known) particle physics
- Schroedinger Equation is 1st order: $i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2}$ 'Take the complex conjugate' (but see Dressel et al. PRL 119 220507 (2017) "Arrow of Time for Continuous Quantum Measurements") Microscopic physics was thought to be symmetric in t

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Simulating \mathcal{T} in decays?

time

We've constructed an angular correlation, a scalar observable, by a dot product of two vectors

 $1+\hat{p}\cdot\hat{J}$

which is odd under *P* as we need

(
$$ec{m{
m p}}$$
 is even, $ec{m{J}}=ec{m{r}} imesec{m{
m p}}$ is odd)

But \vec{J} is odd under T, not even

So we need at least 3 vectors to have a T-odd scalar observable, the scalar triple product $\vec{v1} \cdot (\vec{v2} \times \vec{v3})$

An example \rightarrow

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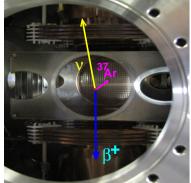
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RIUMF *T* correlation of 3 of 4 momenta

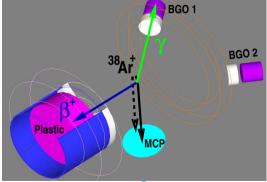
$$\mathbf{t} \to -\mathbf{t} \Rightarrow \vec{p} \propto \frac{\mathrm{d}\vec{r}}{\mathrm{d}t} \to -\vec{p}$$

but $\vec{p}_{\mathrm{recoil}} \cdot \vec{p}_{\beta} \times \vec{p}_{\nu} \equiv \mathbf{0}$

time



$$ec{m{
ho}_{
u}}\cdotec{m{
ho}_{eta}} imesec{m{
ho}_{\gamma}}=-ec{m{
ho}_{ ext{recoil}}}\cdotec{m{
ho}_{eta}} imesec{m{
ho}_{\gamma}}$$
 $\stackrel{t o-t}{\longrightarrow}ec{m{
ho}_{ ext{recoil}}}\cdotec{m{
ho}_{eta}} imesec{m{
ho}_{\gamma}}$



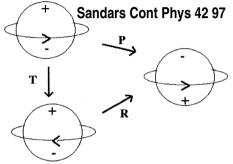
We can test symmetry of apparatus with coincident pairs ☺
 Not exact. Outgoing particles interact → fake X

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EDM in a fundamental particle breaks T: this is exact

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- Landau, Nucl. Phys. 3 (1957) p. 127 Electric Dipole moment $\vec{d} = \sum q_i \vec{r_i}$
- Since the angular momentum is the only vector in the problem, $\vec{d} = a\vec{J}$
- Under T, $\vec{J} \stackrel{t \to -t}{\to} -\vec{J} \quad \vec{d} \stackrel{t \to -t}{\to} +\vec{d}$
- If the physics is invariant under *T*, this is a contradiction, $\Rightarrow a = 0$

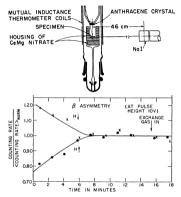


The other logical possibility: there are 2 states, with opposite sign of the EDM, and *T* just formally changes one state to the other.
For most fundamental particles, we know there aren't 2 states
Why do we know the electron doesn't have 2 states?
E.g. some polar molecules have a dipole moment listed in tables, which produces degenerate states and does not break *T* ...]

parity

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Parity broken, why not *T***ime?**



Immediately after P arity was seen to be totally broken in β decay (' ν left-handed') Wu, Ambler, Hayward, Hopper, Hobson, PR 105 (1957) 1413 Many T-odd observables were proposed:

PHYSICAL REVIEW

VOLUME 106, NUMBER 3

Possible Tests of Time Reversal Invariance in Beta Decay

J. D. JACKSON,^{*} S. B. TREIMAN, AND H. W. WYLD, JR. Palmer Physical Laboratory, Princeton University, Princeton, New Jersey (Received January 28, 1957)

Need scalar triple products of 3 vectors: observables involving spin

 $D\hat{J} \cdot \frac{\vec{p_{\beta}}}{E_{\beta}} \times \frac{\vec{p_{\nu}}}{E_{\beta}}$ $R\vec{\sigma}_{\beta} \cdot \hat{J} \times \frac{\vec{p}_{\beta}}{E_{\beta}}$ TRINAT has *D* ideas are consistent with $\mathcal{T} < 0.001$ (We're looking for \mathcal{T} that could still be big:) but some has been found \rightarrow

12/32

parity

time

mtv

🖻 in atoms

PHYSICAL REVIEW

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Possible Tests of Time Reversal Invariance in Beta Decay

J. D. JACKSON,* S. B. TREIMAN, AND H. W. WYLD, JR. Palmer Physical Laboratory, Princeton University, Princeton, New Jersey (Received January 28, 1957)

 \mathcal{CP} discovered in $K\bar{K}$ meson decays in 1963, though not much (Cronin and Fitch Nobel prize 1980) Quark eigenstates in the weak interaction: To explain some weak decays. $|u\rangle \rightarrow |d\rangle + \epsilon |s\rangle$ i.e. $|u\rangle \rightarrow \cos(\theta_{c})|d\rangle + \sin(\theta_{c})|s\rangle$ For 3 families of particles, \rightarrow 3x3 unitary "CKM" matrix between $|d\rangle$, $|s\rangle$, $|b\rangle$ There is one complex phase, which leads to this type of CP A reason for 3 generations of particles?

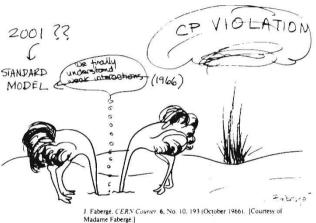
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That one phase is consistent with \mathcal{CP} in $K\bar{K}$ and $B\bar{B}$ systems

There have been hints in $K\bar{K}$ and $B\bar{B}$ of more CP than in the standard model,

time

 $p\bar{p} \rightarrow \mu^+\mu^+ \text{ or } \mu^-\mu^- CP$ at 3.6 σ Abazov PRD 2014 Fermilab; so this 2001 cartoon was a little premature \rightarrow



T2K ν_{μ} oscillations different from $\bar{\nu_{\mu}}$ at 2 to 3 σ Nature 580 339 (2020) *CP* could have some utility for cosmology \rightarrow Sakharov JETP Lett 5 24 (1967) used CP to generate the universe's excess of matter over antimatter:

- CP,
- baryon nonconservation, and
- nonequilibrium.

But known CP is too small by 10¹⁰, so 'we' need more to exist. Caveats:

- You could use CPT [Dolgov Phys Rep 222 (1992) 309]
- \bullet We need CP in the early universe, not necessarily now

So we look for more \mathcal{CP} . How is this related to \mathcal{T} ?

that seems a little abstract concrete demonstrative example (Ramsey-Musolf at INT 2020 \mathcal{CP} explaining T2K's ν vs. $\bar{\nu}$

T2K's ν vs. $\bar{\nu}$ result lets heavy N decay this way in some models

P in atoms time extras --- 🖸 🔂 🔍 Search .washington.edu/talks/WorkShops/int_20_2b/People/Ramsev-Musolf_M/Ramsev-Musolf.pdf - + 150% Neutrinos and the Origin of Matter Heavy neutrinos decay out of equilibrium • in early universe Majorana neutrinos can decay to particles and antiparticles Rates can be slightly different (CP violation) • $\Gamma(N \to \ell H) \neq \Gamma(N \to \bar{\ell} H^*)$

 Resulting excess of leptons over anti-leptons partially converted into excess of quarks over anti-quarks by Standard Model sphalerons

parity	

${\ensuremath{\mathcal{T}}}$ is related to CP by the "CPT Theorem"

"All local Lorentz invariant QFT's are invariant under CPT" Schwinger Phys Rev 82 914 (1951)

Lüders, Pauli, Bell 1954

● Gravity → not flat:

K meson experiments Adler PhysLettB 364 (1995) 239 test *CPT* to within 1000x expected from quantum gravity

• Strings not 'local' Proofs still pursued \rightarrow

	Contents lists available at ScienceDirect	The second secon
and the state	Studies in History and Philosophy	Studies in Theory and Parameter of Modern Physics
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ELSEVIER	journal homepage: www.elsevier.com/locate/shpsb	
On the CPT theorer	n	CrossMark
Hilary Greaves 🏘, Teruji	Thomas ^{b,1}	
^a Somerville College, Oxford OX2 6HD, UK ^b Wolfson College, Oxford OX2 6UD, UK		
	ABSTRACT	
^b Wolfson College, Oxford OK2 6UD, UK	We provide a careful development and rigorous proof of the CPT theo mainstream (Lagrangian) quantum field theory. This is in contrast to the axiomatic frameworks, and non-rigorous proof-stetches in the mainstrea CPT transformation for a general field directly, without appealing to the representations, and in a manner that is clarily related to the requirement	usual rigorous proofs in purely m approach. We construct the enumerative classification of hts of our proof. Our approach
^h Wolfson College, Oxford OK2 GUD, UK A R T I C L E I N F O Article history: Received 21 December 2012 Received 10 mexicad form 25 September 2013 Accepted 7 Oxforder 2013	We provide a careful development and rigorous proof of the CPT theo mainstream (Lagrangian) quantum field theory. This is in contrast to the axiomatic frameworks, and non-rigorous proof-stektes in the mainstrea CPT transformation for a general field directly, without appealing to the representations, and in a manner that is clarity related to the requiremen- applies equally in Minkowski spacetimes of any dimension at least thm between classical and quantum field theories: the quantum CPT theorem in The key mathematical tool is that of complexification; this tool is central to but plays no over role in the usual mainstream approaches to CPT.	isual rigorous proofs in purely m approach. We construct the enumerative classification of uts of our proof. Our approach ee, and is in principle neutral as a natural classical analogue

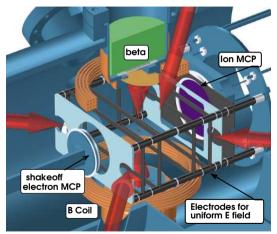
When citing this paper, please use the full journal title Studies in History and Philosophy of Modern Physics

Assuming CPT, $CP \Leftrightarrow T$ in most physics theories The matter excess then motivates T searches trinat

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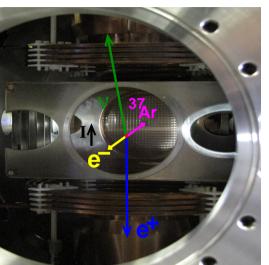
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[®]™³⁷K decay geometry [™]



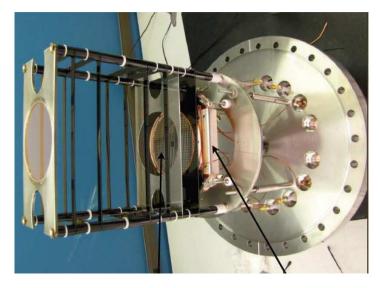
- β , recoil nucleus
- shakeoff e^- for TOF trigger The decay pattern shown on the right is helicity-forbidden if the ν goes straight up



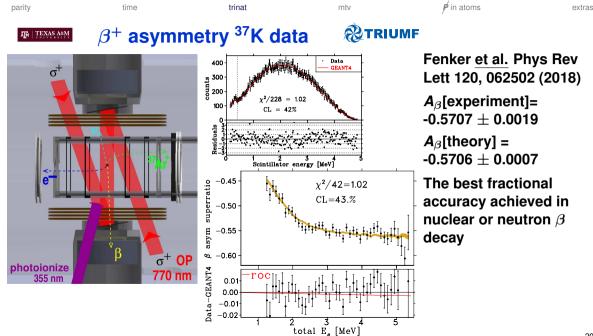


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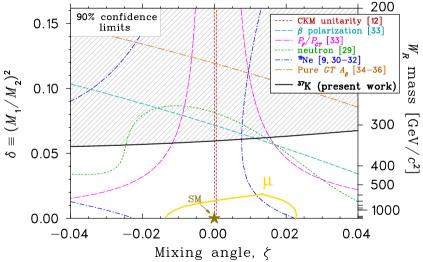
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14 inch CF flange **Electrostatic field** delay-line anode for position info no stray wires Low-Z (glassy carbon, titanium) to minimize β^+ scattering







Extra W' with heavier mass, couples to wrong-handed ν_B LHC *M*'_W > 3.7 TeV 90%

p in atoms

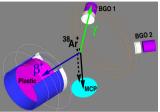
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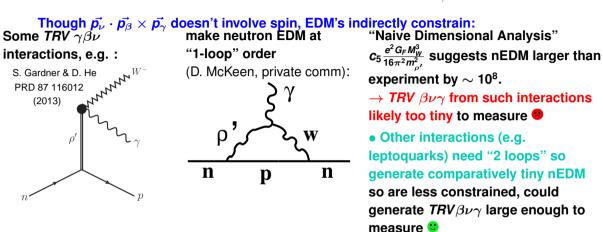
$\mathfrak{B}^{\mathsf{TRIUMF}}$ The nucleon: a special place for γ 's Harvey Hill Hill PRL 99 261601 (2007); EFT with SM interactions combined in the nucleon: goal was extra γ production by medium-energy ν 's **QCD** Weak E&M $\mathcal{L} = \frac{-4c_5}{m_{\rm nucleon}^2} \frac{eG_F V_{ud}}{\sqrt{2}} \epsilon^{\sigma \mu \nu \rho} \bar{p} \gamma_{\sigma} n \bar{\psi}_{eL} \gamma_{\mu} \psi_{\nu L} F_{\nu \rho}$ Gardner, He PRD 2013: looked for contributions to radiative n decay. Noticed QCD antisymmetry led to a scalar triple product of momenta ⁽²⁾: $|\mathcal{M}_{c5}|^2 \propto rac{lm(c_5 g_V)}{M^2} rac{E_e}{D_e k} (ec{p_e} imes ec{k_\gamma}) \cdot ec{p_ u}$ Needs non-SM QCD-like physics, scale $M \sim 10$'s of MeV Particles strongly interacting with themselves but weakly interacting with us are also possible dark matter candidates See the 'SIMP miracle' Hochberg et al. arXiv:1402.5143



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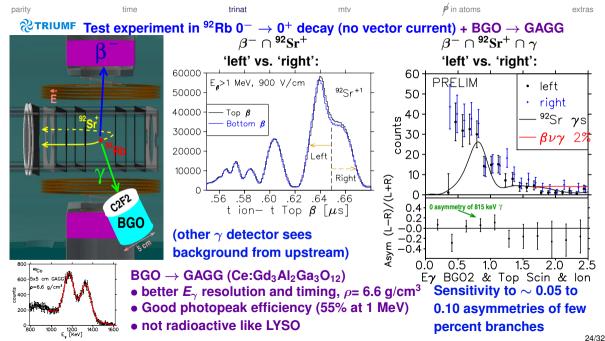
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Other constraints

• Direct constraint from $n \rightarrow p \ \beta \nu \gamma$ branch $\propto |c_5|^2$ Bales PRL 2016: $3.4 \pm 0.2 \times 10^{-3}$ (theory 3.1×10^{-3}) $\Rightarrow \frac{\text{Im}(c_5)}{M^2} \leq 8MeV^{-2} \Rightarrow {}^{37}\text{K} TRV$ asym can be ~ 1

23/32



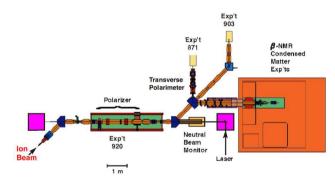
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©TRIUMF Laser-Polarized beam at TRIUMF/ISAC

C.D.P. Levy et al. / Nuclear Physics A 746 (2004) 206c-209c



• 50-70% polarization, 20-50% efficient Re-stripped +1 beam deliverable to several beamlines

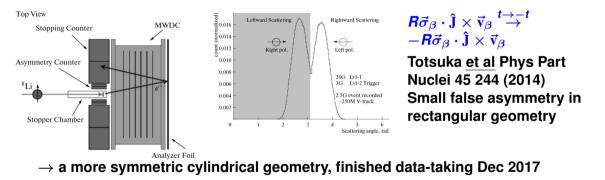
• Used for aligned ²⁰Na β correlation (2nd-class current comparison with ²⁰F) K. Minamisono PRC 84 055501 (2011)

⁸Li *R*, Jiro Murata, Rikkyo U.

TRV possibilities include $^{36}{\rm K}$ *E* and $^{20}{\rm Na}$ β -delayed α energy shift (Clifford PRL 50 (1983) 23)

parity	time	trinat	mtv	🖻 in atoms

Mott scattering Time reversal Violation progress



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mtv

P in atoms is tiny



time

A "chargeless" weak interaction: the biggest Standard Model prediction 1973 ν scattering Gargamelle CERN





Atoms are mostly chemistry, i.e. electromagnetism, which respects Parity Some *P* effects are ppb

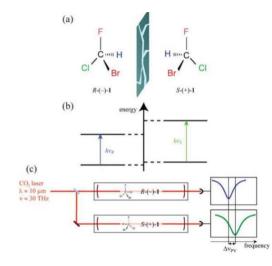
Consider 2 signatures: Molecular enantiomer energies; Controlled Stark-*P*[´]

interference

parity	time	trinat	mtv

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Molecule binding energy depends on handedness



Letokhov, Difference of energy-levels of left and right molecules due to weak interactions. Phys Lett A (1975) 53 275

Darquie <u>et al.</u> CHIRALITY 22 870 (2010) Progress Toward the 1st Observation of P in Chiral Molecules by High-Resolution Laser Spec

 \bullet Very small $\sim 10^{-16}$ to 10^{-14} energy shifts.

Astrobiology 18 (2018) Selection of Amino Acid Chirality via ν Interactions with ¹⁴N in $\vec{E} \times \vec{B}$ Fields M.A. Famiano, R. N. Boyd...

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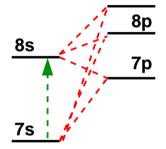
🖻 in atoms

extras

Forbidden transition in an atom

Consider an E1 "electric dipole" transition between S states.

- This is $\not P$ and forbidden.
- (M1 "magnetic dipole" transition is allowed... I avoid a lot of neat physics by setting it to zero.)
- Populate the 8S state, wait for it to decay by a single blue photon to the 7S. Rate \sim interaction², lifetime \sim 10¹⁰ sec (300 yr). With 10⁹ trapped atoms, you get about 6 per minute \rightarrow .
- 9 months for ppt accuracy (a million events).
- Instead we want something linear in the weak interaction:

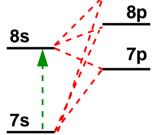


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Stark - Weak interference: flip E field

This idea is credited to the Bouchiat's (theorist-experimentalist couple) Wieman in Boulder made the best 'chargeless weak interaction" measurement in Cesium atoms in 1999. Emulating Boulder Cs scheme in Fr: $|A_{7s \rightarrow 8s}|^2 = |E1_{Stark} + E1_{PNC}|^2$ $\approx |E1_{stark}|^2 + 2E1_{Stark}E1_{PNC}$

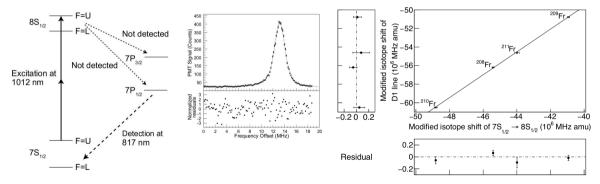


The interference term is $\sim 10^{-9}$ of an allowed E1 transition amplitude (rather than 10^{-18}). By picking an E field one can make the asymmetry about 10^{-3}

• Then calculate (or, preferably, measure) $E1_{Stark}$ to extract $E1_{PNC}$



M. Kalita <u>et al.</u> Phys Rev A 97 042507 (2018) Field shift(8s)/(7s) 1.228 \pm 0.019 exp. vs. 1.234 \pm 0.010 atomic theory



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extras

Discrete symmetries in β decay

time

- Parity *P* symmetry
 How to test *P* symmetry experimentally
 Only left-handed ν so far
- There is time-reversal symmetry violation *T* in nature

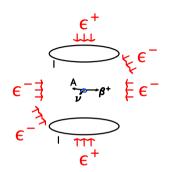
There should? be more: 'baryogenesis'

- TRIUMF Neutral Atom trap (me) β decay
- Laser-polarized beamline at TRIUMF

[Non-decays: • v mixing Nature 580 339 (2020)

- *P* in Fr atoms; Searches for electric dipole moment;
- Using ultra-cold neutrons;
- hopefully atomic fountain of francium atoms]

https://trinat.triumf.ca/talks/triumf-summer-undergrad-talk-2021-jb



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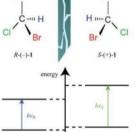
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Preview: Weak interaction breaks parity: Consequences?

'Pulsar kicks'

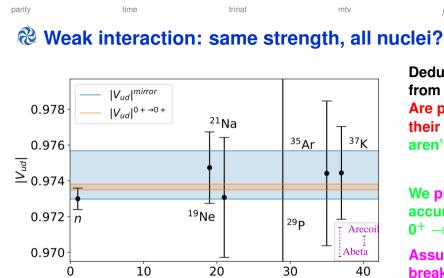


Fuller PRD 2003 Forced $p + e^- \rightarrow n + \nu$ $W(\theta) = 1 + \frac{\langle m_l \rangle}{l} A_{\nu} \cos(\theta_{\hat{1}})$ B field polarizes p's Need ν_e to include 10^{-8} admixture of $m_{\nu} \sim \text{keV}$ Earthling's amino acids are all left-handed



Letokhov PLA'75 Darquie CHIRALITY 2010 $\Delta E \sim 10^{14-16} \text{eV}$ Not Enough for left-handed bugs to win, so \rightarrow Spin-polarized SN ν 's could preferentially zap wrong-handed amino acids Finding the right environment for spin-polarized amino acids? e.g. : Astrobiology 18 (2018) Selection of Amino Acid Chirality via ν Interactions with ¹⁴N in $\vec{E} \times \vec{B}$ Fields M.A. Famiano, R.N. Bovd (TRIUMF EEC 90's)...

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Deduced V_{ud} from mirror decays Are people overestimating their uncertainties? We aren't ⁽¹⁾

TEXAS A&M

p in atoms

We project to reach 0.0005 accuracy, as good as any $0^+ \rightarrow 0^+$ except ^{26m}Al.

Assumes 5% isospin breaking calculation.

Hayen and Severijns, arXiv:1906.09870 (June 2019)

A of initial state

Simulations: E_{γ} signature and backgrounds

• Classical bremsstrahlung $\propto 1/E_{\gamma}$

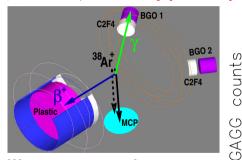
time

 Any time-reversal violating interaction involves β , ν and $\gamma \Rightarrow$ 4-body phase space $\propto E_{\gamma}(Q - E_{\gamma})^3$ Bernard PLB 593 (2004)

counts

trinat

mtv

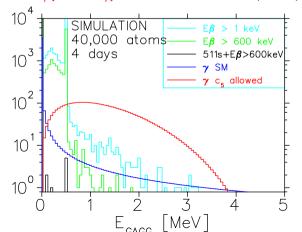


We are concentrating on:

• $E_{\gamma} > 511 \text{ keV}$

parity

• the β^+ in the opposite detector



P in atoms

arity	time

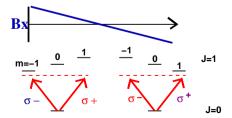
trinat

mtv

in atoms

extras

Magneto-optical trap: perturb atoms Zeeman Optical Trap (MOT) Raab et al. PRL 59 2631 (1987) ε 3 σ^+ σ^{-} **Damped harmonic** 8 oscillator $\varepsilon = \hat{s} \cdot \hat{k}$ **e**+3

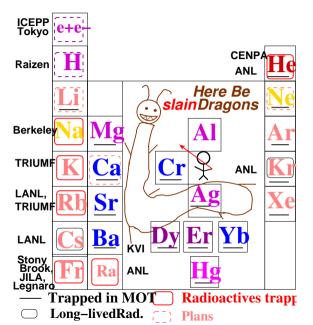




trinat

mtv

What elements can be laser cooled?



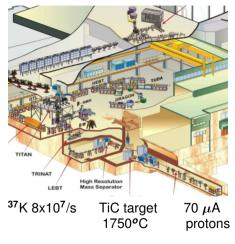
parity	time	trinat

mtv

🖻 in atoms

extras

®TRIUMF TRIumf Neutral Atom trap at ISAC



main TRIUMF cyclotron 'world's largest' 500 MeV H⁻ (0.5 Tesla)



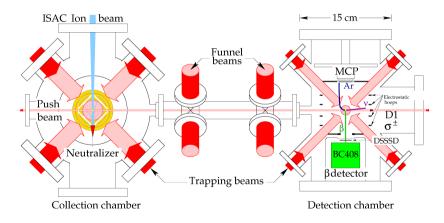
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mtv

RIUMF TRINAT plan view

- Isotope/Isomer selective Avoid untrapped atom background with 2nd trap
- 75% transfer

• 0.7 mm cloud for β -Ar⁺ $\rightarrow \nu$ momentum



 \bullet Spin-polarized 99.1 $\pm0.1\%$

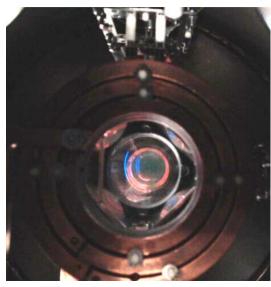
parity

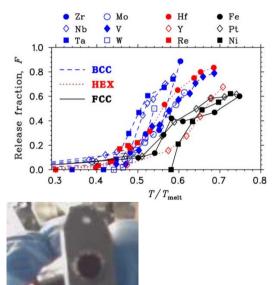
time

trinat

mtv

Neutralizer and Collection trap

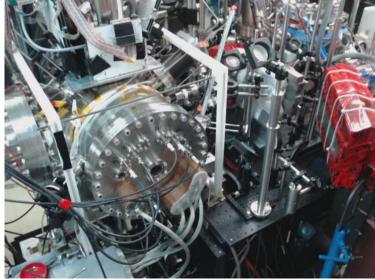


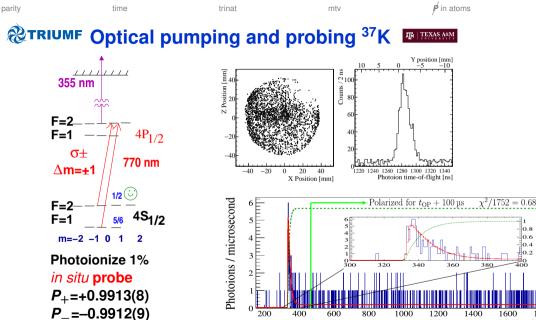


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🖻 in atoms

©TRIUMF TRINAT lab: "tabletop experiment"





Fenker NJP 2016

0.2 1200 1400 1600 1800 Time since MOT off [µs]

0.8

0.6 0.4

0.2

Polarization

0.6

04