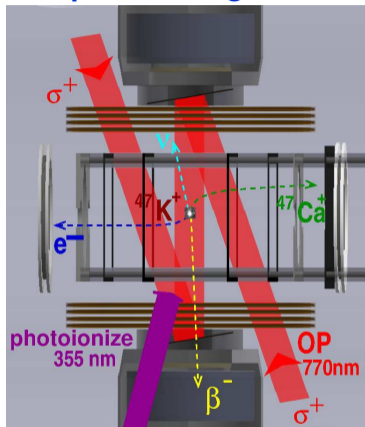


## Isospin breaking and time reversal symmetry in $^{47}\text{K}$ beta decay



- get  $\nu$  momentum from the decay products
- Spin-polarize  $^{37}\text{K}$   $99.1 \pm 0.1\%$  by direct optical pumping

- Isospin symmetry and “isobaric analog states”
- Sensitivity to time-reversal breaking enhanced in isospin-forbidden  $\beta$  decay  $^{47}\text{K}$
- $^{47}\text{K}$  isospin breaking experiment: preliminary results



A. Gorelov  
**B. Kootte\***  
 J.A. Behr



J. McNeil  
 Undergrad:  
 H. Gallop,  
 Waterloo  
 C. Luktuke,  
 Waterloo



UNIVERSITY  
 OF MANITOBA  
 G. Gwinner



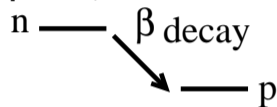
D. Melconian  
 J. Klimo  
 M. Vargas-Calderon

Supported by NSERC, NRC through TRIUMF, DOE, RBC Foundation

\*co-spokesperson

## Isospin: The neutron and proton are isospin projections of the isospin-1/2 “nucleon”. One consequence of isospin symmetry:

neutron  $\beta^-$  decay is to the proton, its **isobaric analog**



tritium  $\beta^-$  decay is also to its **isobaric analog**

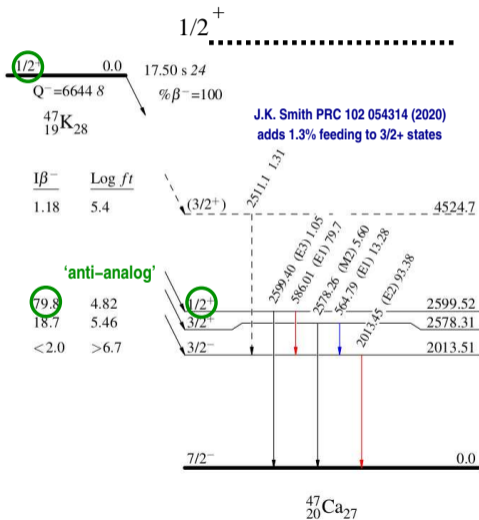


Contributions from Fermi operator  $\tau^\pm$

(Only changes n to p)

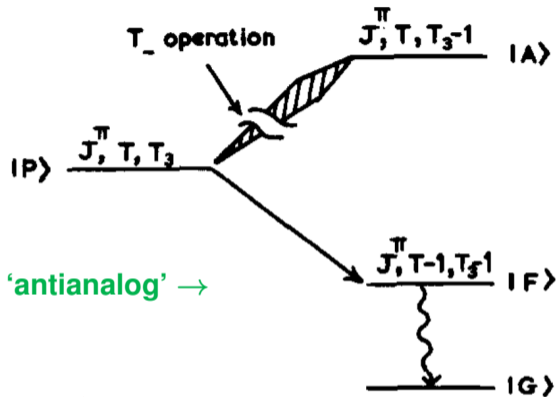
and Gamow-Teller  $\sigma \cdot \tau$

(Can flip spin and isospin)



$^{47}\text{K}$  decay to its isobaric analog is energetically forbidden, so is purely Gamow-Teller, unless isospin mixing of analog and “antianalog” configurations lets Fermi contribute.

Fermi/Gamow-Teller interference changes  $\beta$  decay angular correlations that we measure.

$\mathcal{T}$  in isospin-hindered  $\beta^-$  decay Barroso and Blin-Stoyle, PL 45B 178 (1973)'antianalog'  $\rightarrow$ **Any  $\mathcal{T}$  decay experiment should answer:**

- Does interaction between outgoing particles mimic  $\mathcal{T}$ ? (We hope we can reach the  $D < 10^{-3}$  level of such false  $\mathcal{T}$ )
- Have null EDM's ruled you out? (Not if we reach  $D < 10^{-2}$ )

$$D \hat{\mathbf{J}} \cdot \frac{\vec{p}_\beta}{E_\beta} \times \frac{\vec{p}_\nu}{E_\beta} \xrightarrow{t \rightarrow -t} -D \hat{\mathbf{J}} \cdot \frac{\vec{p}_\beta}{E_\beta} \times \frac{\vec{p}_\nu}{E_\beta}$$

$$D = \sqrt{\frac{J}{J+1}} y / (1 + y^2) \sin(\alpha_V - \alpha_A)$$

$$\text{with } y = \frac{|M_F|}{|M_{GT}|}$$

In this system,  $\sin \alpha_V = -i \frac{\langle F | V_{\mathcal{T}} | A \rangle}{\langle F | V_{\text{Coul}} | A \rangle}$

So for  $\mathcal{T}$  physics mixing antianalog  $|F\rangle$  with analog  $|A\rangle$ , then  $V_{\mathcal{T}}$  is only competing with  $V_{\text{Coul}}$ , not  $V_{\text{strong}}$ ,

enhancing  $\alpha_V$  by  $\sim 10^2$  or  $10^3$  😊

- Has your experiment been done better? (Our goal is 3x better than Calaprice et al.  $^{56}\text{Co}$ , and complementary to NOPTREX neutron scattering resonances for parity-even isospin-breaking interactions)

# Measuring Analog-antianalog mixing for its own sake

N. Auerbach, B.M. Loc arXiv:2101.06199v3

$\bar{A}\bar{A}$  mixing explains isospin-forbidden particle decays,  $\Gamma_A$ , where  $A$  is a well-defined single resonance.

$\bar{A}$  configuration is typically part of several eigenstates:

HO estimate:  $\langle \bar{A} | V_C | A \rangle = 0.35 \frac{\sqrt{n_1 n_2}}{2T} \frac{Z}{A^{2/3}} \text{ MeV}$

$^{88}\text{Sr}$  250 keV Skyrme interactions agree: 250 to 310

	HO	Experiment
$^{71}\text{As}$	300	$28 \pm 4$ Severijns PRC 71 064310 2005 <b>Fragmented <math>\bar{A}</math></b>
$^{56}\text{Co}$	160	$2.9 \pm 0.5$ Markey PRC 26 287R 1982 <b>Fragmented <math>\bar{A}</math></b>
$^{47}\text{K}$	190	<b><math>\bar{A}</math> might be one state</b> 😊

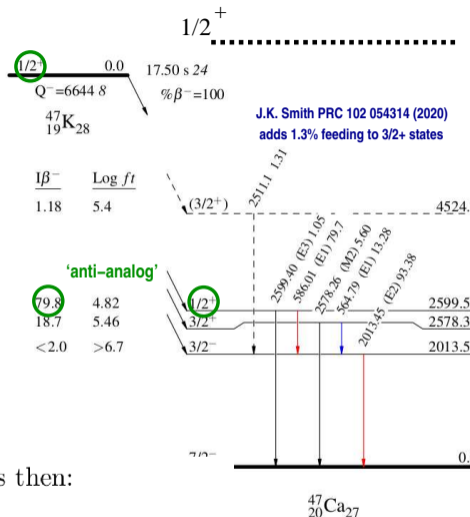
$^{47}\text{K}$ ,  $^{47}\text{Ca}$  are near shell closures 20 and 28 so structure is simpler

The analog is:

The anti-analog  $|\bar{A}\rangle$  is then:

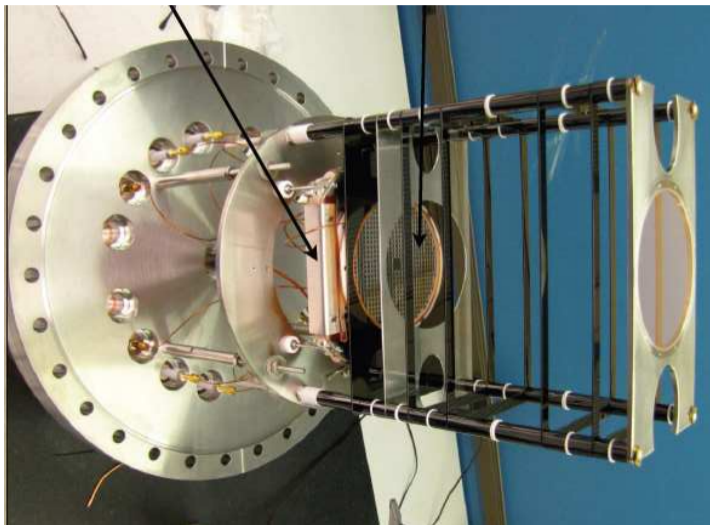
$$|A\rangle = \frac{1}{\sqrt{2T}} \left[ \sqrt{n_1} |j_1^{n_1-1}(n) j_1(p) j_2^{n_2}(n)\rangle + \sqrt{n_2} |j_1^{n_1}(n) j_2^{n_2-1}(n) j_2(p)\rangle \right]$$

$$|\bar{A}\rangle = \frac{1}{\sqrt{2T}} \left[ \sqrt{n_2} |j_1^{n_1-1}(n) j_1(p) j_2^{n_2}(n)\rangle - \sqrt{n_1} |j_1^{n_1}(n) j_2^{n_2-1}(n) j_2(p)\rangle \right].$$





## ion MCP assembly



**14 inch CF flange**

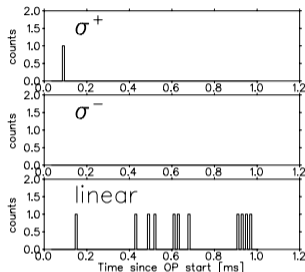
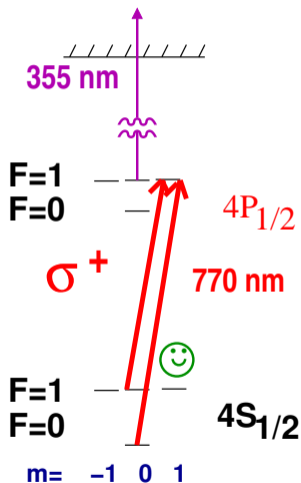
**Electrostatic field**

**delay-line anode for  
position info**

**no stray wires**

**Low-Z (glassy carbon,  
titanium) to minimize  $\beta^+$   
scattering**

## Optical pumping of $I=1/2$ $^{47}\text{K}$



(tight cuts on timing wrt pulse laser and center position exclude background:  
H. Gallop. U. Waterloo)

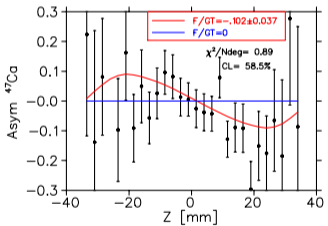
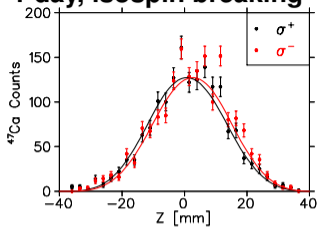
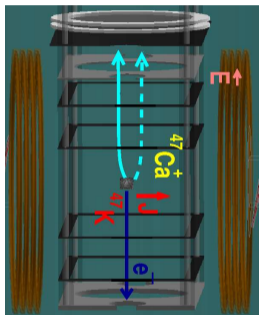
We alternate trap/optical pumping  
Apply circularly polarized light along  $z$  quantization axis.

Once we start OP cycle, atoms increase spin to maximum, then stop absorbing  
If light is linearly polarized, atoms keep absorbing.

When excited, a pulsed laser has enough energy/photon to photoionize (a small fraction) of them.

11 photoions while linearly polarized,  
1 photon circularly polarized  $\rightarrow$   
nuclear polarization  $96 \pm 4\%$

# 1000 atoms trapped for 1 day, isospin breaking <sup>47</sup>K 29 Jan without weak mag



- Nonzero <sup>47</sup>Ca asymmetry wrt spin ⇒ a nonzero  $M_{\text{Fermi}}$
- $M_F / M_{GT} = -0.102 \pm 0.037 \text{ stat} \pm 0.018 \text{ syst} \Rightarrow \langle \bar{A} | V_{\text{Coulomb}} | A \rangle = -78 \pm 28 \text{ stat} \pm 14 \text{ syst keV}$
- $A_{\text{recoil}}$  is damped at extreme Z by a  $\sim 6\%$  bkg from untrapped <sup>47</sup>K, measured by dedicated 'poof' tests
- Apparatus is symmetric: X projection flat at  $1\sigma$  to 0.05; Unpolarized data has X, Z projections flat  $\sim 0.01$
- $\beta'$ 's fire the eMCP with  $\sim 20\%$  quantum efficiency– these we measure to be  $\sim 0.002$  correction

To do: check TOF spectrum to refine 2+ 3+ 4+ ratio

Weak magnetism correction



## $\beta$ -recoil (Confirms sign of polarization)

Fit to similar numerical integration, including pointlike  $\beta$  detector and a 2 MeV photon.  
Scaling with number of +,-; DSSSD XY strips; solid angle from Z shift of trap; 0.99 for  $\langle \cos \theta_\beta \rangle$ ; divide asymmetry by 1.023 as in 37K to account for backscatter; all put into calculation, not data

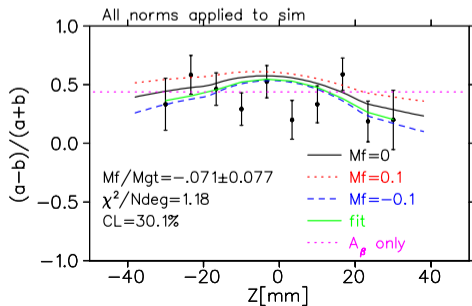
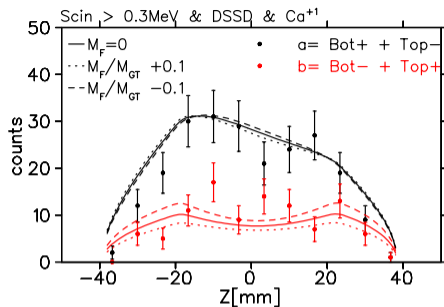
There is an evident change in the  $\beta$  &  $^{47}\text{Ca}$  asymmetry with radius; acceptance in the other dimension with trap offset + the  $\gamma$ -ray momentum forces some  $\text{Ca}^{+1}$  to miss.

$M_f/M_{gt} = -0.071 \pm 0.077 \text{ stat} \pm 0.033 \text{ syst}$

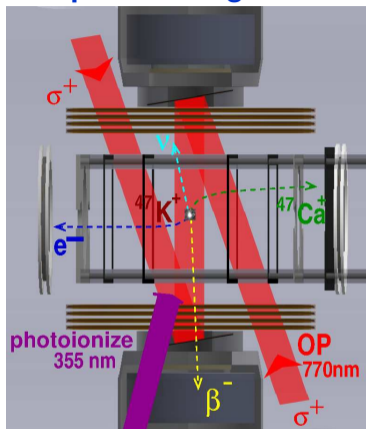
Uncertainties: polarization  $0.96 \pm 0.04 \rightarrow$  uncertainty 0.022

Scale Z by 0.9 to make the distribution fit better by eye, uncertainty  $\rightarrow 0.025$  ("E field uncertainty")  
assume 20% uncertainty backscatter  $\rightarrow 0.0024$   
Added in quadrature  $\rightarrow M_f/M_{gt} = -0.071 \pm 0.084$

To do: Quantify weak magnetism correction  $\sim 0.01$



## Isospin breaking and time reversal symmetry in $^{47}\text{K}$ beta decay



$p_{^{37}\text{Ar}}$ : uniform  $\vec{E}$ ,  
 MCP for TOF and position  
 $p_{\beta}$ : from  $\delta E + E$   
 $\rightarrow p_{\nu}$  event-by-event  
 Spin-polarized  $^{47}\text{K}$   $96 \pm 4\%$

- Preliminary results suggest a nonzero measurable Fermi component to the main branch of  $^{47}\text{K}$   $\beta$  decay  
 We're measuring something that isn't 'zero' 😊!  
 Our standard model tests have all agreed with predictions 😞
- The possibly large value of the Coulomb-induced isospin breaking is predicted because near-closed shell  $^{47}\text{Ca}$  has only one bound state with same  $J^{\pi}$  as  $^{47}\text{K}$ 
  - Measuring *isospin* in  $^{47}\text{K}$  will determine sensitivity to parity-even *isospin*  $\mathcal{T}$  interactions via future  $D\vec{I} \cdot \vec{v}_{\beta} \times \vec{v}_{\nu}$

**Isospin: The neutron and proton are isospin projections of the isospin-1/2 “nucleon”**

**It provides another degree of freedom for antisymmetrization of fermion wf's under exchange of identical fermions**

**Isospin is an abstract symmetry,**

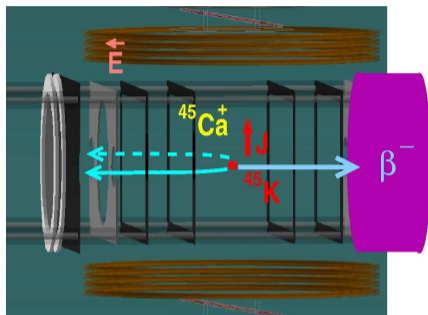
**yet Wigner's  $SU(4) = SU(2)$  spin  $\times$   $SU(2)$  isospin explains most quantized states in light nuclei: one can classify complexity of states in light nuclei by the number of  $SU(4)$  configurations (Ormand and Vogel)**

**The Coulomb interaction breaks isospin. We will interpret our measurement of 47K isospin breaking in terms of a Coulomb matrix element.**

**QCD only breaks isospin a little, because with  $m_u \neq m_d$ . More is commonly invoked phenomenologically to explain the Nolen-Schiffer anomaly of mirror nuclear masses: this has consequences for the isospin breaking needed for 1% corrections to the absolute strength of  $\beta$  decay and for our other  $\beta$  decay project**



## $D \vec{l} \cdot \vec{v}_\beta \times \vec{v}_\nu$ in atom trap: Features, Systematics



- Collect recoils going into 4 pi with electric field of 1 kV/cm
- Full reconstruction of recoil and beta momenta
- Point source: we know where it is (by sampling photoionization) and it doesn't move when we flip the polarization

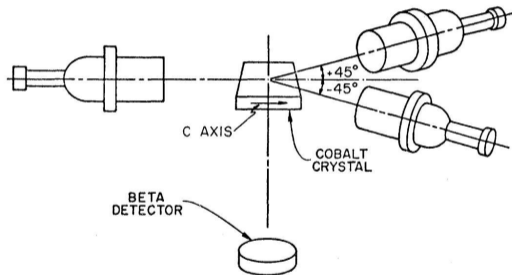
$D$  Uncertainties / 100 scaling from Melconian PLB 649 270 (2007)

	$B_\nu$	Improvements	Projected
Cloud position $\sigma^\pm$	1.3	$\pm 500 \mu\text{m} \rightarrow \pm 20 \mu\text{m}$	0.05
Cloud size/Temp	0.3	" "	0.03
MCP Position cal	1.0	DLA+ mask	$\leq 0.1$
$\hat{x}$ -OP alignment	0.25	Geometry is $\perp$	$\leq 0.02$
E field	0.2		$\leq 0.1$

- Any stray polarization along wrong axis is deadly, a lowest-order fake  $D$ : Measure with singles asymmetry for recoils and  $\beta$ 's

## $^{56}\text{Co}$ $\mathcal{T}$ experiment

Asymmetry of the  $45^\circ$   $\gamma$  detectors with nuclear alignment



“Test of time-reversal invariance in the beta decay of  $^{56}\text{Co}$ ”

Calaprice, Freedman, (Princeton);  
Osgood, Thomlinson (BNL)  
PRC 15 381 (1977)

$$E_1 = -0.01 \pm 0.02$$

$\log(ft) = 8.7$ , yet known allowed:  
 $E_\beta$  spectrum, no  $\beta$ - $\gamma$  correlation)

$y = -0.13 \pm 0.02$  PRC 26 287R (1982)  
Markey, Boehm (RIP Felix 2021)

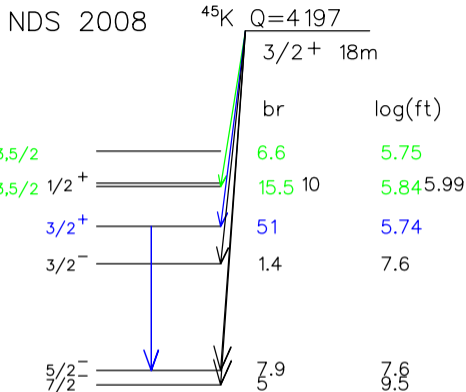
$$V_{\text{Coul}} = 2.9 \text{ keV}, V_{\mathcal{T}} = 54 \pm 110 \text{ eV}$$

(J.L. Mortara Ph.D. thesis 1999 UCB

$$E_1 = -0.001 \pm 0.006$$

$$\Rightarrow V_{\mathcal{T}} = 5 \pm 33 \text{ eV}$$

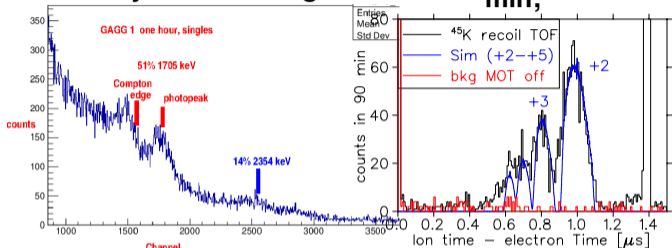
We believe we can measure  $D$  in  $^{47,45}\text{K}$   
much more accurately than  $E$  in  $^{56}\text{Co}$ ,  
but we must find a case with  $|M_{GT}|$ ,  
 $V_{\text{Coul}}$ , and  $\mathcal{T}$  N-N matrix elements to  
allow complementary or better  
sensitivity to  $V_{\mathcal{T}}$



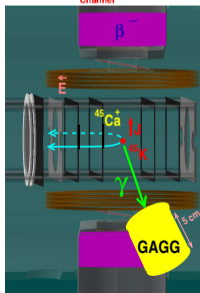
**51% branch to  $3/2^+$  state in  $^{45}\text{Ca}$ .  
Should include the antianalog configuration,  
 $\langle F | V_{\text{coul}} | A \rangle \sim 5$  to  $50$  keV ?**

**$A_{\beta}$ ,  $A_{\text{recoil}}$  would answer 15.5% branch to  $1/2^+$   $3/2^+$   $5/2^+$ ?**

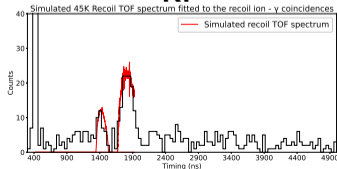
**$^{45}\text{K}$  decay to antianalog**



**shakeoff  $e^-$  & recoil  
clean even for  $t_{1/2}=18$   
min;**



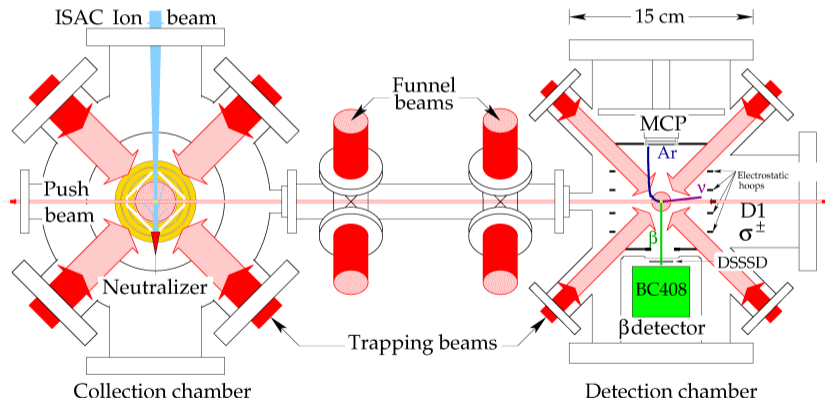
**$\gamma$  & recoil is a challenge  
that will be cleaner in  
 $^{47}\text{K}$ :**





# TRINAT plan view

- Isotope/Isomer selective
- 75% transfer
- Avoid untrapped atom background with 2nd trap
- 0.7 mm cloud for  $\beta$ -Ar<sup>+</sup>  $\rightarrow$   $\nu$  momentum



- Spin-polarized  $99.1 \pm 0.1\%$

# Neutralizer and Collection trap

