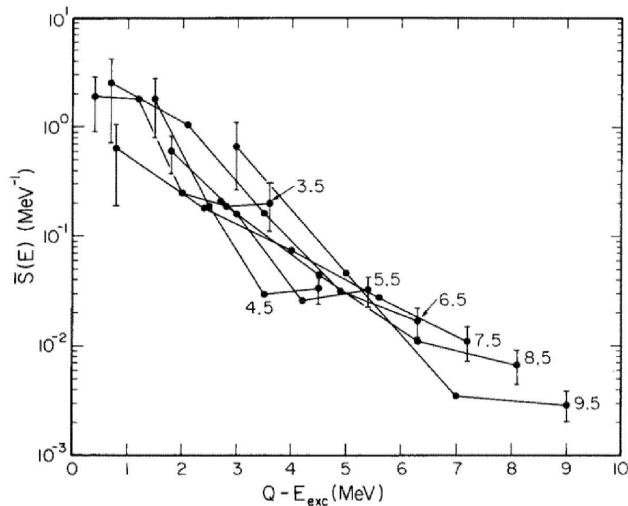


# $\beta$ -decay Phenomenology of Nuclear Fission Products

*J. A. Behr, P. Vogel /  $\beta$ -decay phenomenology*



**Behr and Vogel**

**My 1982 summer research**

**Nucl Phys A411 199 (1983)**

**Summed database 100 fission products**

**Hints that nuclear giant resonances are fed**

**Higher energy  $\nu$ 's dominated by a few cases with large energy release**

[trinat.triumf.ca/publications-1/reactor\\_nus\\_2021.pdf](http://trinat.triumf.ca/publications-1/reactor_nus_2021.pdf)



## Reactor $\nu$ energies with an atom trap

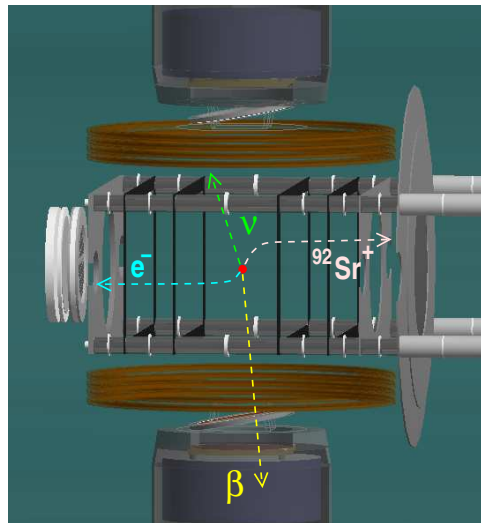
- Nuclear reactors produce a lot of  $\nu$ 's
  - (Total # measured)/(calculated) =  $0.92 \pm 0.04$   
A remarkable success– but what is missing?
  - Discrepancy is worse between 5 and 7 MeV
- How well is the source understood?
- We measure the energy spectrum of  $\nu$ 's produced by a particular type of  $\beta$  decay, to test theory understanding

### TRIUMF Neutral Atom Trap (TRINAT)

How atom traps work

How we can measure  $\nu$  energy

First results for  $^{92}\text{Rb}$  decay



# TRIUMF Neutral Atom Trap collaboration:



D. Melconian



A. Gorelov

J.A. Behr

Undergrad  
D. Prins UBC



J. McNeil



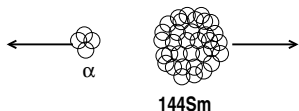
UNIVERSITY  
OF MANITOBA

**M. Anholm**

G. Gwinner

Support: NSERC, NRC through TRIUMF, DOE, State of Texas

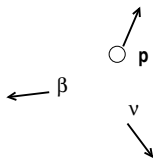
## $\nu$ was invented to solve an experimental puzzle



$^{144}\text{Sm}$

$$p_{\alpha} = p_{^{144}\text{Sm}}$$

$$E_{\alpha} = 3.183 \text{ MeV, always}$$



“Controversy and Consensus: Nuclear  $\beta$  decay 1911-1934” Springer 2000, eds. Hiebert, Knobloch, Scholz (C. Jensen)

$\beta$  decay: A continuous  $E_e$  spectrum, not a discrete peak!

Meitner and Hahn 1911, Danysz 1913, experimentally resolved:

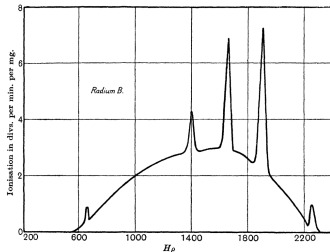


Figure 3.12: The beta spectrum of radium B, obtained by Chadwick and Ellis when they repeated Chadwick's experiment of 1914. Source: Chadwick and Ellis, “Preliminary Investigation” (note 82), p. 277.

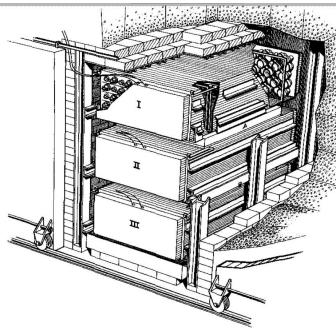
- 1923 Ellis+Wooster: statistical energy conservation
- 1929 Niels Bohr: non-conservation of energy (!?) sought to power stars...?
- 1930 Pauli postulated a new particle (??!!)

How to test?

Probability to interact in a detector follows from the neutron decay rate (Bethe and Peierls, Nature **133** 532 (1934); Robson Phys Rev **83** 349 (1951))

Pauli: “I have done a terrible thing... postulated a particle that cannot be detected.”

# Reactor $\nu$ 's: first direct confirmation by "Inverse $\beta$ decay"



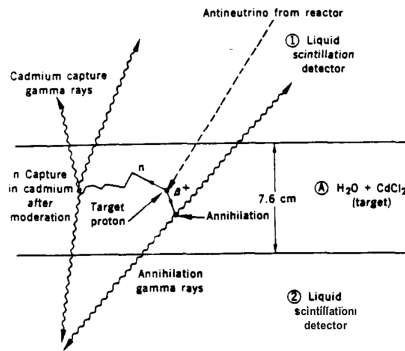
sketch of the equipment used at Savannah River. The

200 liters  
 $4 \times 10^{-6}$  SuperK's

## 1995 Nobel Prize

Nobel Lecture 1995

Fredrick Reines



compared to the expected<sup>2</sup>

With parity violation (1957) prediction is 2x bigger :)

$$\bar{\sigma}_{exp} = (12^{+7}_{-4}) \times 10^{-44} \text{ cm}^2$$

$$\bar{\sigma}_{th} = (5 \pm 1) \times 10^{-44} \text{ cm}^2$$

1st plan: put a detector next to a **nuclear bomb**

Pulsed source, get above natural backgrounds ☺

Must calibrate detector well before experiment ☹

**Reactor** worked better:  
 1956 Science **124** 103  
 C. Cowan, F. Reines, Harrison, Kruse, McGuire (Los Alamos)  
 They thought they could predict the number to  $\sim 30\% \rightarrow$

## TRIUMF Two reactor $\nu$ 'anomalies'

- total  $\nu$  flux is  $92 \pm 4\%$  of expected  $\rightarrow$  extra  $\nu$ ?

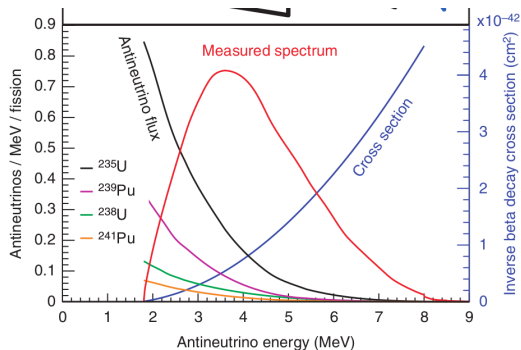
Update Hayes et al. PRL 120 022503 (2018):

Fuel composition changes with time

Still room for  $\sim 5\%$  discrepancy and a sterile  $\nu$

Flux[distance] measurements (PROSPECT) may clarify this

- Disagreement between detectors and computation at  $\nu$  energy 5-7 MeV



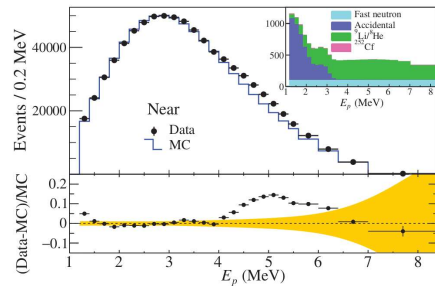
P. Vogel, L.J. Wen, C. Zhang, Nature Comm 6 6935 (2015)

The 5-7 MeV  $\nu$ 's are a fair fraction of the detected  $\nu$ 's



## TRIUMF The 'bump' is now well-measured, but not explained

- Experimental excess over models  $E_\nu$  5-7 MeV  
Seen in reactor experiments RENO and Daya Bay  
Still consistent with PROSPECT first result
- Understanding is needed for ambitious neutrino hierarchy measurement with reactor  $\nu$  oscillations (thus a near detector planned for JUNO)
- There are models with 'new physics' to explain the 'bump' (Barryman, Brdar, Huber PRD 99 055045)
- Nuclear theory generally is now estimating larger uncertainties for weak magnetism and 1st-forbidden decays, making the bump more consistent with less precise theory.
- Nearly half of these 5-7 MeV  $\nu$ 's come from  $0^- \rightarrow 0^+$  decays



RENO PRL 121 201801 (2018)

# TRIUMF Reactor $\nu$ 'anomalies' and $^{92}\text{Rb}$ decay

$^{92}\text{Rb}$   $Q=8104$

$^{92}\text{Rb} \sim 10\%$  of reactor  $\nu$ 's 5-7 MeV

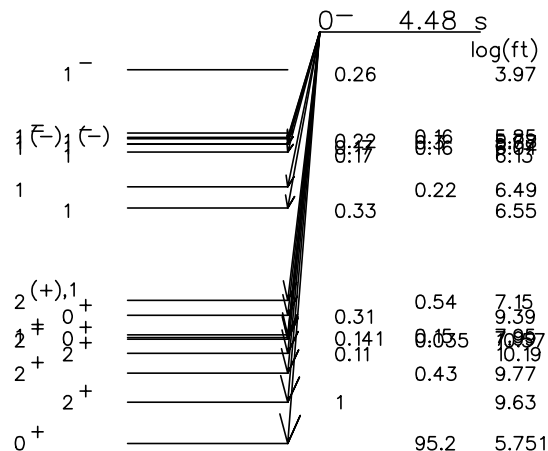
2012 NDS compilation: g.s.  $\rightarrow$  g.s.  
branch  **$95.2 \pm 0.7\%$** , based on Lhersonneau  
et al. PRC 74 017308 (2006) feeding of  
first  $2^+ \rightarrow 0^+ \gamma$   **$3.2 \pm 4\%$**

• Total absorption spectrometer results:  
Zakari-Issoufou et al. PRL 115 102503  
(2015)  **$87.5 \pm 2.5\%$**

Rasco et al. PRL 117 092501 (2016)  
 **$91 \pm 3\%$**

• Conventional thick scintillator + Ge  
done at ANL is in between (E. McCutchan,  
Apr 2018 APS)

**Branching ratio likely under control**





**Is the energy spectrum of  $\nu$ 's predictable?**

**$\beta$  energy spectrum for allowed decay:**

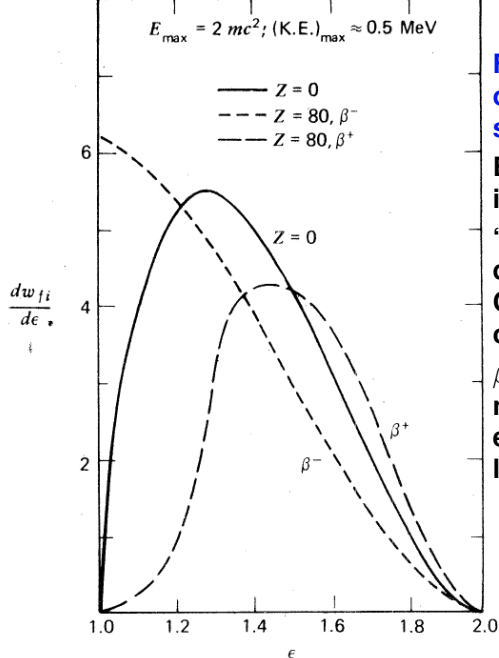
Assuming the outgoing leptons are plane waves,  
Integrate  $p_e^2 dp_e p_\nu^2 dp_\nu \delta(Q - K_e - E_\nu)$  over  $p_\nu$ ,

$$W(p_e) dp_e = \frac{1}{2\pi^3 \hbar^3 c^3} \sum_{\mu} \sum_{m_f} |\langle \mathbf{J}_f m_f r | \mathbf{O}_{\lambda\mu}(\beta) | \mathbf{J}_i m_i r' \rangle|^2 F(Z, K_e) p_e^2 (Q - K_e) \sqrt{(Q - K_e)^2 - m_\nu^2} dp_e$$

with  $\mathbf{O}_{\lambda\mu} = \sum_{j=1}^A (\mathbf{G}_V \tau_{\pm}(j) + \mathbf{G}_A \vec{\sigma}(j) \tau_{\pm}(j))$ . Differentiating  $E^2 = p^2 + m^2 \Rightarrow p dp = E dE$ ,

$$W(E_e) dE_e \propto F(Z, E_e) E_e p_e (E_0 - E_e) \sqrt{(E_0 - E_e)^2 - m_\nu^2} dE_e$$

**(Most forbidden decay operators produce large changes in this energy spectrum)**



## Fermi function effect on $\beta$ energy spectrum

But the outgoing  $\beta$  is not a plane wave

“Coulomb function” distorted by the Coulomb potential of the nucleus.

$\beta^-$  is ‘pulled into’ nucleus... a big effect for high  $Z$  and low  $E_\beta$

## Selection rules

**Fermi**

**G-T**

$\gamma_5$  dominates  $0^- \rightarrow 0^+$

$\sigma \cdot r$  suppressed by  $r/\lambda$   
 but that  $r$  dependence  
 distorts the  $E_\nu$  spectrum

'1st forb. unique'  $2^\pm \leftrightarrow 0^\mp$

One operator  $\Rightarrow$  calculable  
 correlations from a.m.

TABLE I. Allowed and first-forbidden nuclear matrix elements and their selection rules ( $K$  designates the rank of the transition operator, when regarded as a tensor).

| Matrix element                                   | $K$ | $\Delta J$  | $\Delta\pi$ |
|--|-----|---|-------------|
| Allowed $C_V f 1$                                | 0   | 0   | +1          |
| $C_A f \boldsymbol{\sigma}$                      | 1   | 0, $\pm 1$ (no $0 \rightarrow 0$ )  | +1          |
| First for-<br>bidden $C_A f \gamma_5$            | 0   | 0   | -1          |
| $C_A f (\boldsymbol{\sigma} \cdot \mathbf{r}/i)$ |     |   |             |
| $C_V f \mathbf{r}i$                              | 1   | 0, $\pm 1$ (no $0 \rightarrow 0$ )  | -1          |
| $C_V f \boldsymbol{\alpha}$                      |     |   |             |
| $C_A f (\boldsymbol{\sigma} \times \mathbf{r})$  |     |   |             |
| $C_A f i B_{ij}$                                 | 2   | 0, $\pm 1$ , $\pm 2$ (no $0 \rightarrow 0$ ,<br>no $1 \rightarrow 0$ , no $0 \rightarrow 1$ ) | -1          |

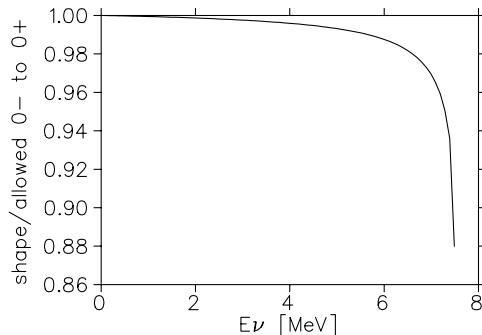


## Non-allowed $\nu$ spectrum from $^{92}\text{Rb}$ decay?

**Corrections beyond Gamow-Teller allowed energy spectra are thought to be important for reactor  $\nu$  spectra**

**Sonzogni, McCutchan, and Hayes PRL 119 112501 (2017) “precisely measured electron spectra for about 50 relevant fission products are needed” to pin down weak magnetism and forbidden correction terms.**

**Historically,  $E_\beta$  spectra in  $0^- \rightarrow 0^+$  decay disagree with theory: worth new technique**



**$0^- \rightarrow 0^+$  correction to allowed  $\beta$  spectrum is theoretically possible [Hayes, Friar, Garvey et al. PRL 2014], removing 0-10%  $\nu$  at highest energy.**



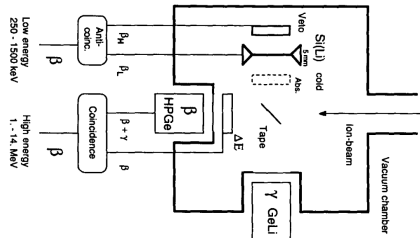
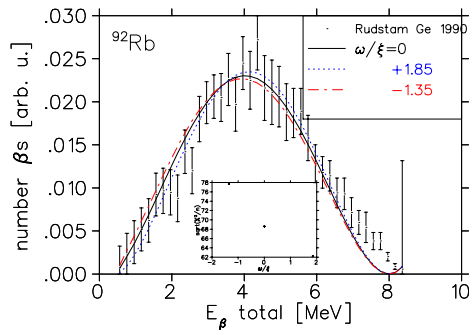
## Other experiments permit $^{92}\text{Rb}$ to be 'non-allowed'

● Reviews typically say  $^{92}\text{Rb}$  beta spectrum is consistent with an allowed shape. This is based on the spectrum from Rudstam et al. ADNDT 45 239 (1990)

Theory expects one matrix element to dominate, producing an allowed spectrum shape

Yet experiments can accommodate the deviations our measurements imply  $\rightarrow$

●  $^{134}\text{Sb}$  with a Paul trap at ANL Siegl et al. PRC 97 035504 (2018): average  $a_{\beta\nu} = 0.47 \pm 0.16$ , attributed to excited state feeding 3%  $\rightarrow$  17%.



O. Tengblad et al. / Integral  $\beta$ -spectra

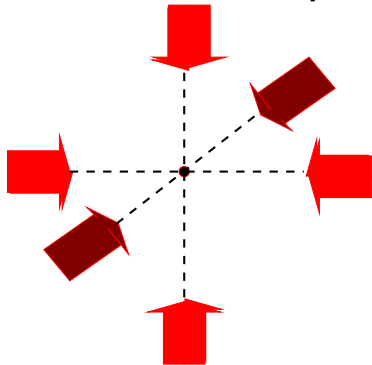


on TRINAT's wall (from co-op H. Norton)



## Magneto-optical trap: damping

For a trap, we want a damped harmonic oscillator  
'Red-detuned' beams provide the "damping"

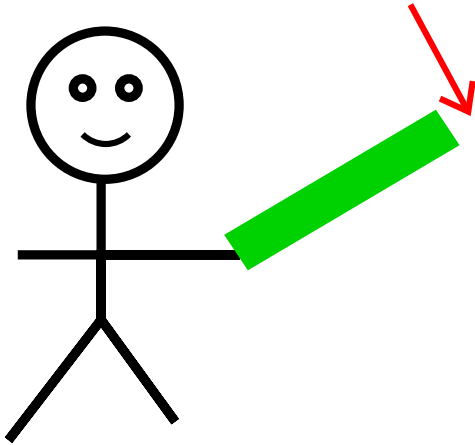


'Optical molasses'

We still need a position-dependent force

# “‘Light sabers’ would make atom traps easy” (H. Norton)

$$\vec{\nabla} \cdot \vec{S} \neq 0$$

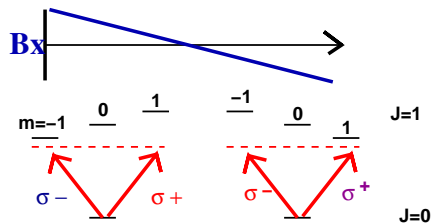
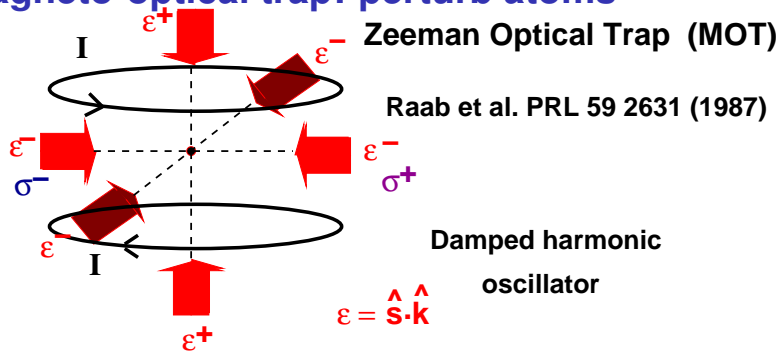


“Optical Earnshaw Theorem”  
(Ashkin + Gordon 1983)

**But light sabers violate Poynting's theorem**



# Magneto-optical trap: perturb atoms



|                                     |        |    |  |    |    |    |  |              |    |
|-------------------------------------|--------|----|--|----|----|----|--|--------------|----|
| ICEPP<br>Tokyo                      | $e+e-$ |    |  |    |    |    |  |              |    |
| Raizen                              | H      |    |  |    |    |    |  | CENPA<br>ANL | He |
|                                     | Li     |    |  |    |    |    |  |              | Ne |
| Berkeley                            | Na     | Mg |  |    | Al |    |  |              | Ar |
| TRIUMF                              | K      | Ca |  | Cr |    |    |  | ANL          | Kr |
| LANL,<br>TRIUMF                     | Rb     | Sr |  | Ag |    |    |  |              | Xe |
| LANL                                | Cs     | Ba |  | Dy | Er | Yb |  |              |    |
| Stony<br>Brook,<br>JILA,<br>Legnaro | Fr     | Ra |  | Hg |    |    |  |              |    |

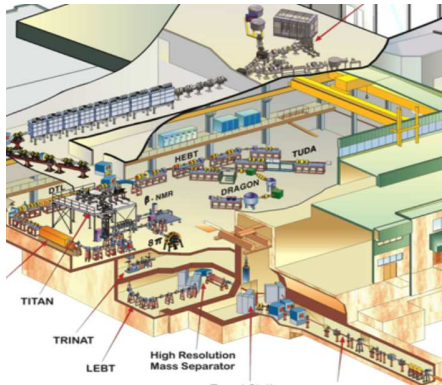
What elements can be laser cooled?

— Trapped in MOT    Radioactives trapped

○ Long-lived Rad.    Plans



# TRiumf Neutral Atom trap at ISAC

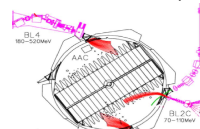


$^{37}\text{K}$   $8 \times 10^7/\text{s}$

TiC target  
 $1750^\circ\text{C}$

$70 \mu\text{A}$   
protons

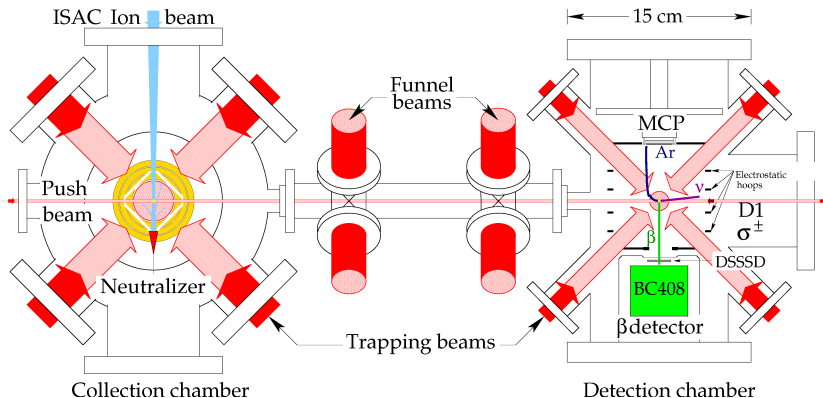
main TRIUMF cyclotron  
'world's largest'  
 $500 \text{ MeV } \text{H}^-$  (0.5 Tesla)





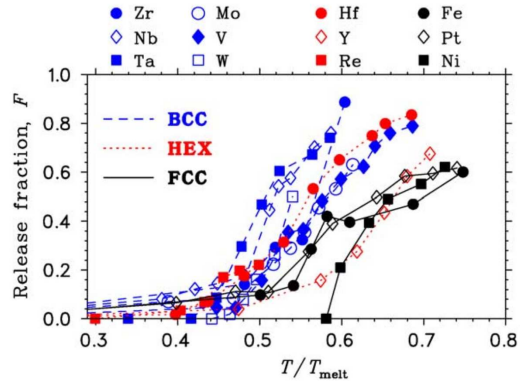
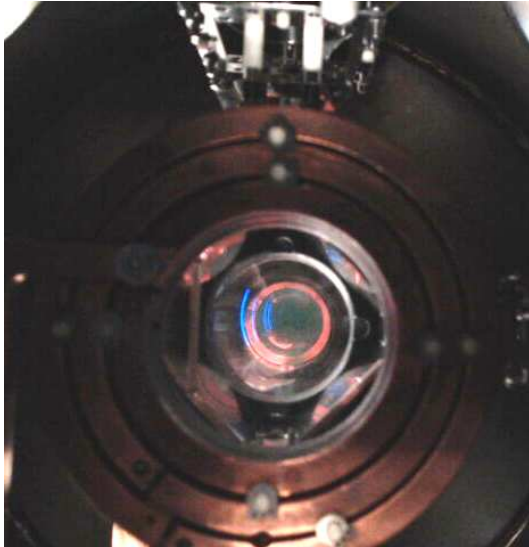
## TRINAT plan view

- Isotope/Isomer selective
- 75% transfer
- Avoid untrapped atom background with 2nd trap
- 0.7 mm cloud for  $\beta$ -Ar<sup>+</sup>  $\rightarrow$   $\nu$  momentum



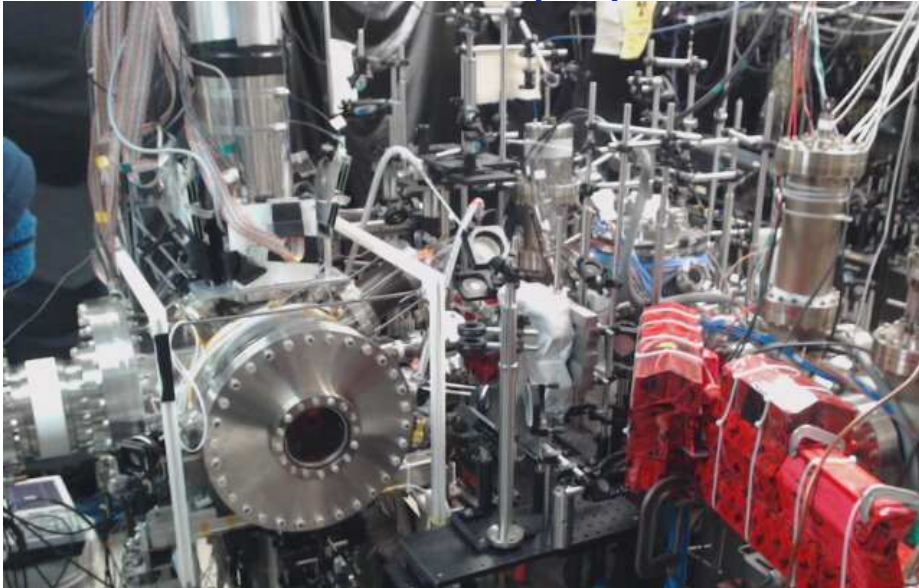
- Spin-polarized  $99.1 \pm 0.1\%$

# Neutralizer and Collection trap



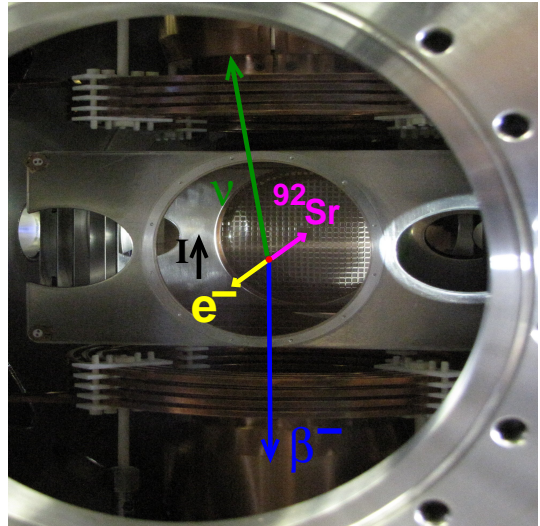
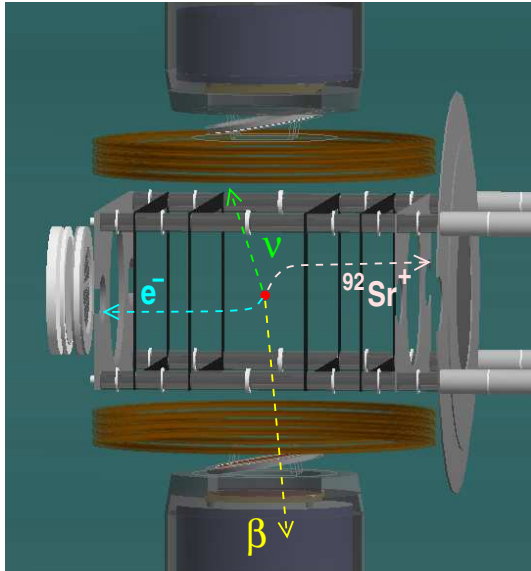


## TRINAT lab: “tabletop experiment”

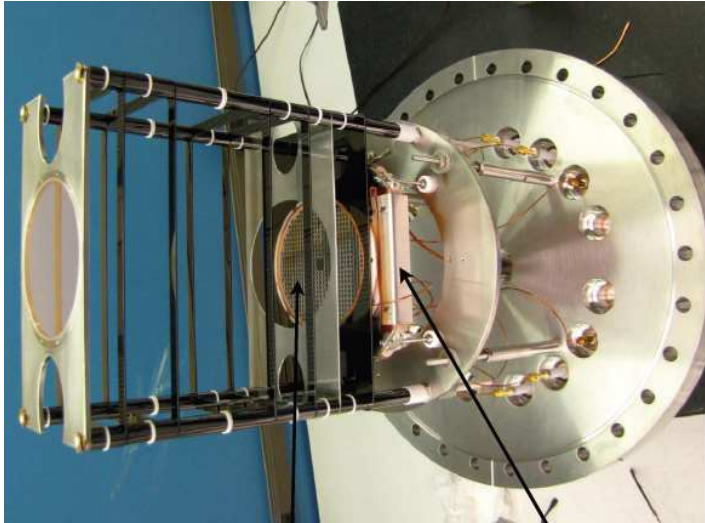




# $^{92}\text{Rb}$ Decay geometry



## ion MCP assembly



**14 inch CF flange**

**Electrostatic field**

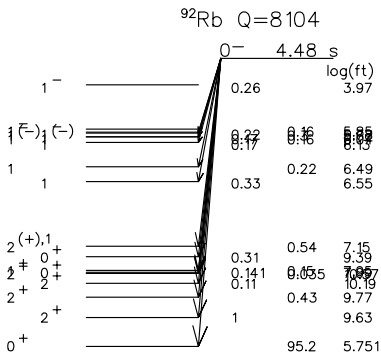
**delay-line anode for  
position info**

**No stray wires**

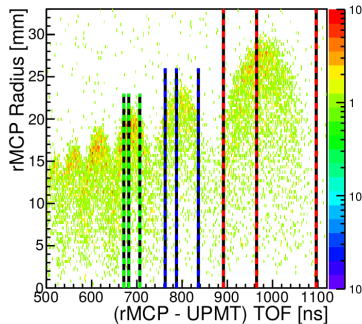
**Low-Z (glassy carbon,  
titanium) and open  
structure to minimize  $\beta^+$   
scattering**



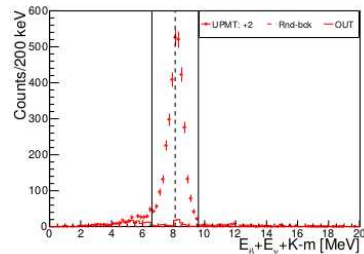
# TRIUMF Data $^{92}\text{Rb}$ (10% of total)



Three experiments (two T.A.S., one careful  $\beta$ - $\gamma$ ) now concur on  $\approx 10\%$  excited state feeding



Scatter plot of recoil TOF vs. recoil radius.



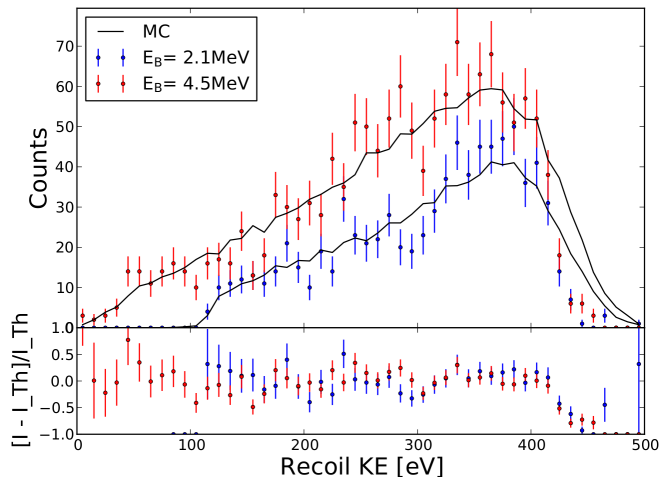
We can separate the decay to the ground state cleanly from the reconstructed total energy



# TRIUMF Preliminary: Reactor $\bar{\nu}$ 's from $0^- \rightarrow 0^+$ $^{92}\text{Rb}$ decay

Determine  $a_{\beta\nu}$  from the recoil energy spectrum:

Warburton PRC 1982:



$$P(E, \theta_{\beta\nu}) = 1 + a \frac{v}{c} \cos(\theta_{\beta\nu})$$

$$a = \frac{1 - \frac{\omega^2}{9\xi_0^2}}{1 + \frac{\omega^2}{9\xi_0^2} - \frac{2\omega m_{\beta}\gamma}{3\xi_0 E_{\beta}}} \xrightarrow{\omega \ll \xi_0?} 1$$

Nuclear matrix elements:

$$\langle i || \sigma \cdot r || f \rangle / R_{\text{nucleus}} = \omega$$

$$\langle i || \gamma_5 || f \rangle \rightarrow \xi_0$$

Which  $\Rightarrow$   $\beta$  spectrum

distorted by:

$$1 + \frac{\omega^2}{9\xi_0^2} - \frac{2\omega m_{\beta}\gamma}{3\xi_0 E_{\beta}}$$

We see that  $a$  changes with  $E_{\beta}$   
 $E_{\nu}$  spectrum changed by  $\lesssim 10\%$



## Reactor $\nu$ energies with your atom trap

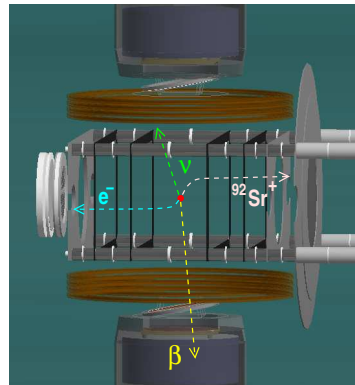
- Nuclear reactors produce a lot of  $\nu$ 's
- Discrepancy between calculation and experiment is worse between 5 and 7 MeV

We measure the energy spectrum of  $\nu$ 's produced by  $0^- \rightarrow 0^+$   $\beta$  decay, to test theory understanding

**TRIUMF Neutral Atom Trap (TRINAT)**

Measured  $\beta$ - $\nu$  correlation to test theory

First results for  $^{92}\text{Rb}$  decay shown at 2020 DNP

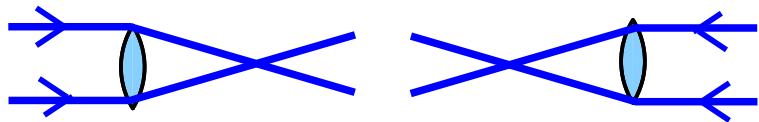


## “Why Optical Traps Can't Work”

Earnshaw Theorem:  $\vec{\nabla} \cdot \vec{E} = 0 \Rightarrow$

no electrostatic potential minimum for charge-free region

“Optical Earnshaw Theorem” (Ashkin + Gordon 1983):



$\Rightarrow$  no 3-D traps from spontaneous light forces  
with static light fields

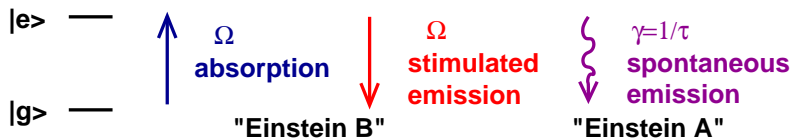
Using Poynting's theorem:

$$\vec{\nabla} \cdot \vec{S} = \frac{c}{4\pi} \vec{\nabla} \cdot (\vec{E} \times \vec{B}) = -\vec{J} \cdot \vec{E} - \frac{\partial u}{\partial t} = 0$$

Dodges !

- Time-dependent forces (pulsed lasers)
- Dipole Force traps (“optical tweezers”)
- Modify internal structure of atom with external fields

## Why atom traps are shallow



$$\frac{dN_g}{dt} = -\Omega N_g + \Omega N_e + \gamma N_e = -\frac{dN_e}{dt}$$

Steady-state  $\Rightarrow =0 \Rightarrow N_e = \frac{\Omega N_g}{\Omega + \gamma}$

Limits:  $N_e \xrightarrow{\Omega \leq \gamma} \frac{\Omega}{\gamma} N_g$  (sure);  $N_e \xrightarrow{\Omega \gg \gamma} N_g$  !!

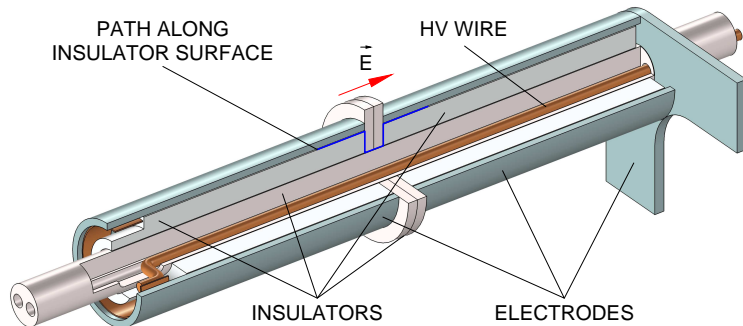
• At high intensity, same # in ground, excited state  
Atomic transition “saturates”

Maximum scattering rate =  $\gamma N_e / N \rightarrow \gamma / 2$

So radiation pressure traps are shallow IF they rely on spontaneous emission

## TRIUMF 'No stray wires'

**Nested insulators: E  
no longer falls across  
dielectric surfaces**



- Argon conditioning
- 1.2 kV/cm reached
- Improved ion MCP mount (as in Hong et al. NIM Seattle-Argonne) in progress
- More compact shakeoff  $e^-$  MCP and wedge-and-strip readout to allow simultaneous ion and  $e^-$  detection.

- Remove  $A_\beta$  background
- Adds  $A_{\text{recoil}}$
- All detectors together for trap diagnostics and for  $\rho$ -independent  $\beta$ -recoil observable