- The most accurate *A*_β measurement Agrees with theory prediction
- Constraints on:

Weak interaction changes within nuclei Non-SM lepton helicities: Left-right symmetric models. 4-fermi contact Lorentz 'scalar', 'tensor'

TRlumf Neutral Atom Trap collaboration:



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intro

- M. Mehlman
- P. Shidling

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M.R. Pearson

ing Undergrad

D. Melconian E. Broatch G. Gwinner S. Supported by NSERC, NRC through TRIUMF, Israel Science Foundation, DOE, State of Texas



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Lepton helicity \rightarrow angular distribution



← This decay pattern needs non-S.M. chirality

I=3/2 \rightarrow I=3/2:

Leptons can't increase nuclear spin any further



One experimental discovery of parity violation

Wu, Ambler, Hayward, Hopper, Hobson, PR 105 (1957) 1413



Wauters 2010 PRC $A_{60Co} = -1.014 \pm 0.020$ [SM -0.987 ± 0.009]

Here A_{β} isn't 1 or -1 or a clean fraction there are 2 operators:

'Fermi' changes n to p

'Gamow-Teller' changes n to p and nucleon spin

au, *Q*, and branch \Rightarrow decay strength $\mathcal{F}t$ We know the Fermi $\mathcal{F}t_0$ from the $0^+ \rightarrow 0^+$ decays, so from $\mathcal{F}t$ we can get the Gamow-Teller strength:



 $\mathcal{F}t$ (Shidling PRC 2014) \Rightarrow

$$\rho = C_A M_{GT} / C_V M_F = 0.5768 \pm 0.0021$$

 $\Rightarrow A_{\beta}[SM] = -0.5706 \pm 0.0007$ main uncertainty is experimental branching ratio

³⁷K isobaric mirror decay: a 'heavy neutron'

A_G physics

 $\Rightarrow A_{\beta}[SM] = -0.5706 \pm 0.0007$ Dominant uncertainty is exp. branching ratio 1st-order recoil-order from E&M moments: Induced tensor $d_1 \approx 0$, Small $\mu \Rightarrow$ small weak magnetism

Recoil-order + Coulomb + finite-size corrections $\Rightarrow \Delta A_{\beta} \approx -0.0028 (E_{\beta}/E_0)$ Holstein RMP 1975



37 K A_B

Isospin mixing contributes 0.0004 uncertainty from shell model

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DFT for isospin mixing has improved its functional Using weighted average for δ_c would $\Rightarrow 0.0004 \rightarrow 0.0005$

intro



³⁷K 8x10⁷/s

TiC target 1750°C

70 µA protons 500 MeV H⁻ (0.5 Tesla)



®TRIUMF TRINAT lab: "tabletop experiment"









- β , recoil nucleus
- shakeoff e⁻ for TOF trigger



This decay pattern is helicity-forbidden if the ν goes straight up, independent of Gamow-Teller/Fermi ratio.







 \bullet 2.8 $\times10^{-3}$ of events in main peak are background from non-trapped atoms

• Conservatively assume polarized between 0 and 100%. \rightarrow A_{β} \times (1.0014 \pm 0.0014)

• These will be removed by MCP position info when we increase to design *E* field

ir	httro 37 K A_{β} A_{β} physical sector $A_{$	ysics	extras	j.a.behr triumf cap17	
		ainties	R TRIUMF		
	Source $\times 10^{-4}$ [†: β scattering]	ΔA_{eta}	A_{eta} = -0.570	7 ±	
	Background (Correction 1.0014)	7	0.0013 (stat)) ±	
	Trap Position	4	0.0012 (syst) ±	
	Trap Sail velocity	5	0.0005 (pol)		
Trap Temperature & width		1	$= -0.5707 \pm 0.0018$		
	BB1 Radius [†]	4			
	BB1 Energy agreement	2	$A_{\beta}[SM] =$		
	BB1 threshold	1	-0.5706 + 0.0007		
	Scintillator threshold	0.3			
	GEANT4 physics list [†]	4	Better relati	ve	
	Shakeoff electon t.o.f. region	3	uncertainty	than	
	SiC mirror thickness [†]	1	¹⁹ Ne –0.0360	0±0.0008	
	Be window thickness [†]	0.9	[Calaprice 1	975]	
	Scintillator or summed [†]	1	and neutron		
	Scintillator calibration	0.1	0.1197+0.00	006	
	Total systematics	12		PRI 2013	
	Statistics	13		20121	
	Polarization	5		2013]	
	lotal uncertainty	18	rnysics $ ightarrow$		



¹Na

.2'O

. 15

+

 $|\mu|$

1²⁴A)

progeny)/A

.25

DN ٧e

.30

.978

.976

.974 > 3

.972

.970

.968

.966

.00

.05

 $(|\mu|)$

- <u>1</u>0

parent

 $0^+ \rightarrow 0^+$ determination of V_{ud} i.e. $\psi[n] \neq \psi[p]$ Salam and Strathdee Nature 1974: phase transitions at very high B fields could drive $V_{ud} \rightarrow 1$ Hardy Towner PLB 1975 applied to the ³⁵Ar A_{β} controversy. ¹⁹Ne Broussard DNP 2016

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Why the weak interaction is 'weak' at low energy

'more massive virtual particles are created for shorter times'

Propagator+vertices: $T \propto \frac{G_{X}(-g^{\mu\nu}+p^{\mu}p^{\nu}/M_{X}^{2})G_{X}}{p^{2}-M_{X}^{2}} \xrightarrow{p << M_{X}} \qquad n \qquad p \qquad X$ $T \propto \frac{G_{X}^{2}}{M_{X}^{2}} \Rightarrow \qquad p \qquad X$ • Decay rates $\propto \frac{G_{X}^{2}G_{X}^{\prime 2}}{M_{X}^{2}}$ or $\propto \frac{G^{2}}{M_{W}^{2}} \frac{G_{X}G_{X}^{\prime 2}}{M_{X}^{2}}$ if process interferes with W (couples to SM-handed ν) e.g. Fierz term $\propto \frac{m}{E_{\beta}}$

• IF $G_X \sim$ electroweak coupling, then 0.1% sensitivity in angular correlations $\rightarrow M_X \sim 6$ or 30 M_W





(or $g_B < 4$, at 2 TeV but LHC7 2 TeV 'bump' had $q\sim 0.5$)





intro



Ave A_{recoil} depends on ρ ; p dependence doesn't





extras



®TRIUMF TRIUMF Neutral Atom Trap: Near Future

We have measured the β asymmetry of ³⁷K decay to be A_{β} =-0.5707 \pm 0.0018

Agrees with theory -0.5706 ± 0.0007 , complements the best β decay measurements

We plan to measure $A_{\beta}[E_{\beta}]$ 3-5 x better, and A_{recoil} with sensitivity to '4-fermion contact' interactions complementary to $\pi \rightarrow e\nu\gamma, \pi \rightarrow e\nu$, and LHC $\rho + \rho \rightarrow e + E_{\perp}$

We also plan a TRV $\beta\nu\gamma$ 3-momentum correlation, first of its type in 1st-generation particles







Re[g_ -0.4 + -0.2 PRELIMINARY -0.05 0.Ò0 0.05 0.0 0.6 0.2 0.4 $Re[g_{a_{1}}^{s} = (C_{s} + C_{s}')/2]$ $(C_{\pi} + C_{\pi}^{\dagger})/C_{\star}$ the Fierz term is 'easier' to constrain but has more competition

-0.05

For scalars coupling to wrong-chirality ν , we compete with our own ^{38m}K β-ν Gorelov 2005

intro 37 K A_{β}		A _B physics	S	extras	j.a.behr triumf cap17					
RIUMF Polarization Improvements RITEXISAN										
SYST $ imes$ 10 ⁻⁴	ΔP		Δ	T						
	σ^{-}	σ^+	σ^{-}	σ^+	• pellicle					
Initial T	3	3	10	8	mirrors:					
Global fit v. ave	2	2	7	6	less β^+					
S ^{out} Uncertainty	1	2	11	5	scattering					
Cloud temp	2	0.5	3	² define 7						
Binning	1	1	4	³ trim B o	iradients					
B _z Uncertainty	0.5	3	2	7						
Initial P	0.1	0.1	0.4	0.4 Improve	e S ₃ flipping					
Require $\mathcal{I}_+ = \mathcal{I}$	0.1	0.1	<u>0.1</u>	0.2 and gradi	ents					
Total SYSTEMATIC	5	5	17	14 ● add flip	ping of <i>B_z</i>					
STATISTICS	7	6	21	17 • higher-	oower					
B. Fenker New J. Phys 18 073028 photoionizing laser										
2016				• gentler	RAC-MOT					
$P(\sigma^+) = +0.9913(8)$ $T(\sigma^+) = -0.9770(22)$ • Uncertainty \propto										
$P(\sigma^{-1}) = -0.9912(9) T(\sigma^{-1}) = -0.9761(27)$ (1-P)										

MSSM and β decay correlations

Profumo, Ramsey-Musolf, Tulin PRD 75 075017 2017 $C_S+C'_S$ can be 0.001 in MSSM in 1-loop order including mixing



FIG. 2. Feynman diagrams relative to supersymmetric contributions giving rise to anomalous amplitudes in β decay processes.

Include mixing of:

• left and right sfermions (this is where β decay can help; constraints are said to be few)

• sfamily mixing (already tightly constrained, e.g. by $\mu \rightarrow e \gamma ...$) Effective 4-fermi scalar and tensor couplings are generated that contribute to **b**_{Fierz} and spin correlation observables like **B**_{ν} as large as 0.001.

Weakly-coupled W' still has electric charge



Does $\sigma e^+ + e^- \rightarrow W^+ + W^-$ double for W'? Depends on the cut for W: typically this cut (explicitly listed in PDG) excludes low-mass W because of serious background

nucleon form factors

Herczeg Prog Part Nucl Phys 46 (2001) 413 pointed out need for form factors

n $\mathbf{g}_{g_s} = \mathbf{g}_{g_s} =$

 \rightarrow (2016) PRD 94 054508

 $g_s = 0.97 \pm 0.12 \pm 0.06, \ g_T = 0.987 \pm 0.051 \pm 0.020$

 g_{s} =1.02 \pm 0.10 Gonzalez-Alonso, Camalich PRL 112 042501 (2014) isospin symmetry



"2nd-class" weak interactions would violate isospin symmetry when quarks are combined by QCD into nucleons. "Induced tensor" *d* is near zero in isobaric mirror decay.

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This result is complementary to other nuclear β decay (Sumikama PRC 2011) in models where 2nd-class currents change with system (Wilkinson EPJA 2000)

Babar set best 3-generation constraints PRL 2009

 $au^-
ightarrow \omega \pi^-
u_ au$

What elements can be laser cooled?



intro

$$\begin{aligned} \mathbf{A}_{\text{obs}}^{\text{SR}}(\mathbf{E}_{e}) &= \frac{1 - s(\mathbf{E}_{e})}{1 + s(\mathbf{E}_{e})} = \mathbf{A}_{\text{obs}} \\ \mathbf{s}(\mathbf{E}_{e}) &= \sqrt{\frac{r_{1}^{-}(\mathbf{E}_{e})r_{2}^{+}(\mathbf{E}_{e})}{r_{1}^{+}(\mathbf{E}_{e})r_{2}^{-}(\mathbf{E}_{e})}} \end{aligned}$$

Gay, T.J. and Dunning, F.B. Rev. Sci. Instrum. 63 (1992) 1635

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B. Plaster et al. PRC 86 (2012) 055501